

Laser-Induced Breakdown Spectroscopy (LIBS)

presented at
*Principles of and Recent Advances in Laser
Micro/Nano Manufacturing Processes*

June 1-4, 2010 - Evanston, IL

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UNIVERSITY

Outline

1. Introduction to LIBS
2. Physics of the plasma formation and observation
3. Instrumentation
4. Advanced techniques

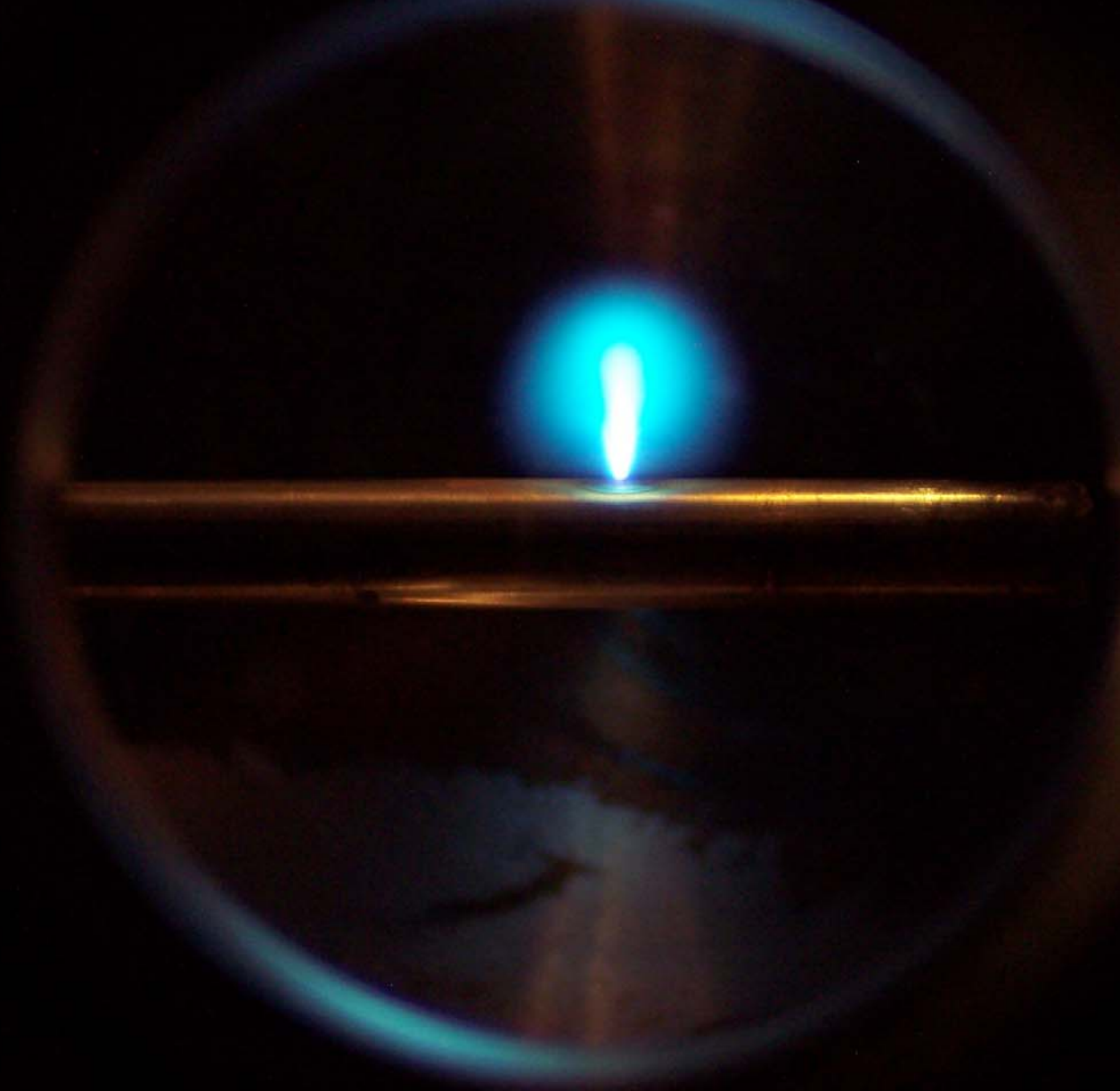


LIBS Defined

One sentence?

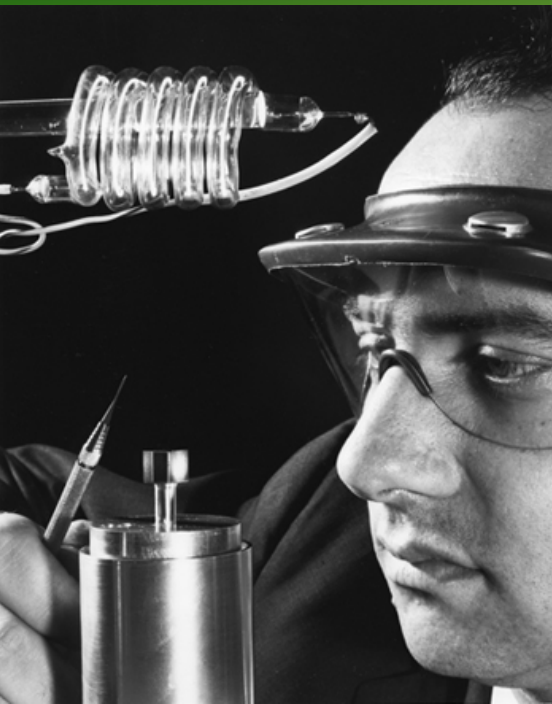
A spectrochemical technique which utilizes an intense laser pulse to determine the atomic/elemental composition of a sample via generation of a high-temperature micro-plasma followed by time-resolved optical spectroscopy.



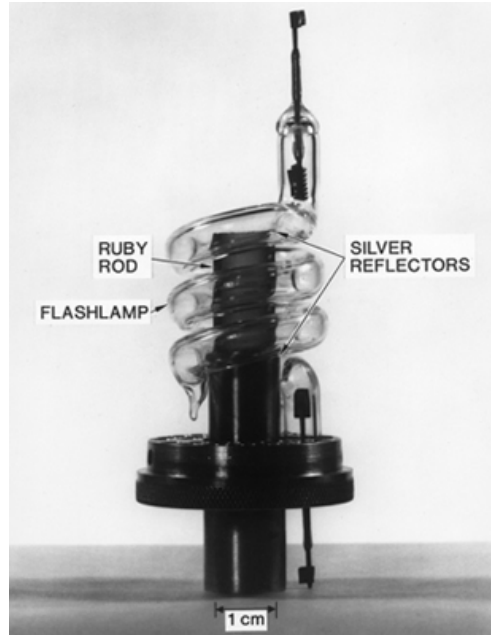


08/01/20

History



1960
Maiman, first ruby laser



1962
Brech, Cross; Birth of LIBS: detection of spectrum from ruby laser induced plasma

1965
Zel'dovich, Raizer, First theoretical model for laser breakdown of a gas

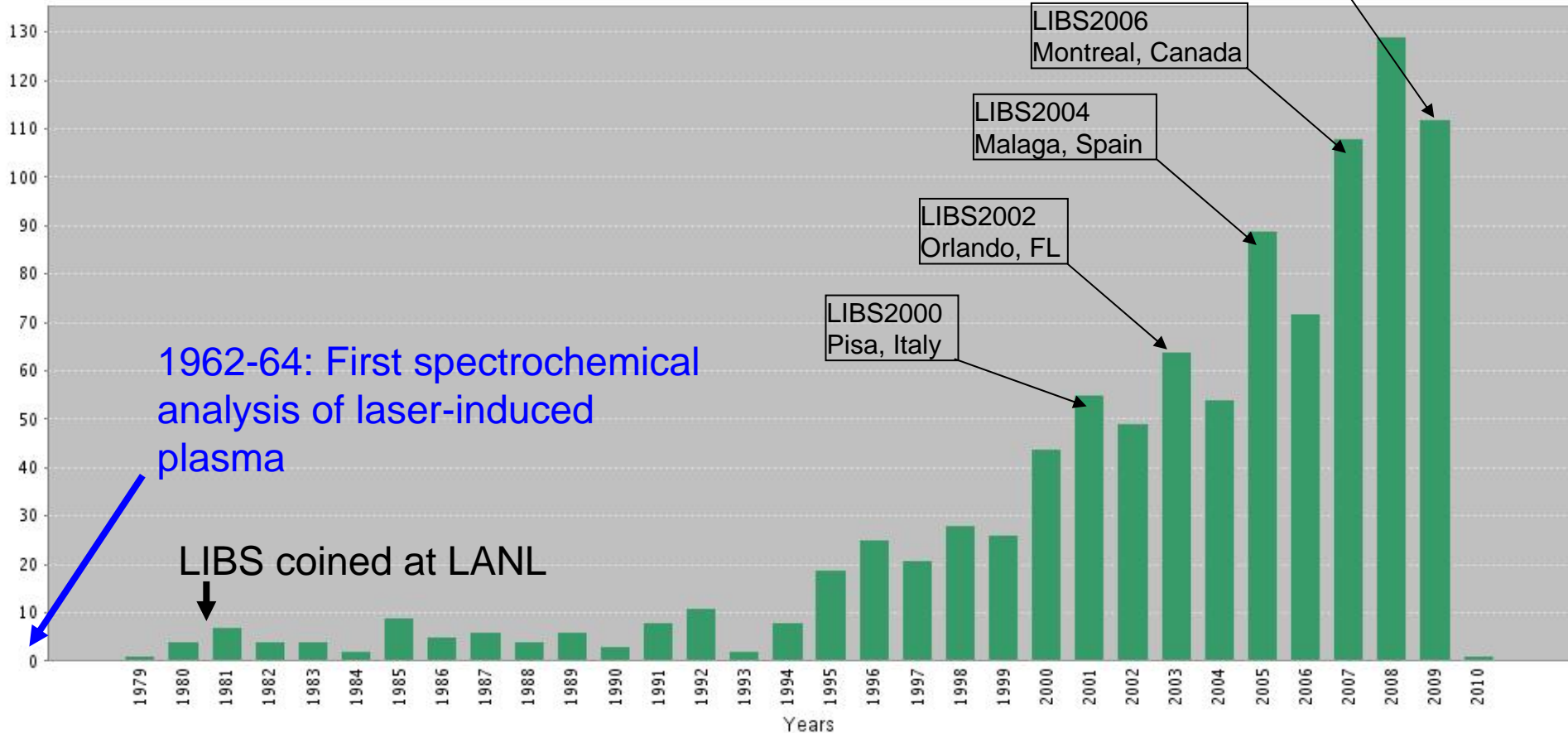
1964
Runger et al. First direct spectrochemical analysis by LIBS













History: “laser-induced breakdown”^{**} @ Web of Science

*Searches for “laser-induced plasma spectroscopy”
would increase hits

Published Items in Each Year



Where we publish

| Field: Source Title | Record Count | % of 1722 | Bar Chart |
|--|--------------|-----------|---|
| SPECTROCHIMICA ACTA PART B-ATOMIC SPECTROSCOPY | 315 | 18.2927 % |  |
| APPLIED SPECTROSCOPY | 156 | 9.0592 % |  |
| APPLIED OPTICS | 100 | 5.8072 % |  |
| JOURNAL OF APPLIED PHYSICS | 68 | 3.9489 % |  |
| JOURNAL OF ANALYTICAL ATOMIC SPECTROMETRY | 60 | 3.4843 % |  |
| APPLIED PHYSICS LETTERS | 46 | 2.6713 % |  |
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| JOURNAL OF PHYSICS D-APPLIED PHYSICS | 32 | 1.8583 % |  |
| APPLIED PHYSICS A-MATERIALS SCIENCE & PROCESSING | 31 | 1.8002 % |  |
| Field: Source Title | Record Count | % of 1722 | Bar Chart |

(303 Source Title value(s) outside display options.)

What's Driving the Interest in LIBS?

- mid-80's: reliable, small, inexpensive lasers
- mid-80's: intensified charge-coupled devices (ICCD)
- 90's – 00's: femtosecond pulsed lasers
- 90's – 00's: broadband spectrometers and Echelle spectrometers
- 00's: microchip lasers
- 00's: portable systems
- 00's: chemometric techniques



LIBS Applications

- **industrial processes**
 - analysis of steam generator tubes in nuclear power stations
 - grading of powdered pellets for glass melts
 - analysis of treated wood in recycling centers
 - grading of iron-ore slurry prior to pelletizing
- **environmental analysis**
 - quantification of heavy metal content in soils, sand, and sludge
 - measurement of lead content in paint
 - water quality assessments
 - hazardous waste remediation
 - atmospheric sampling
- **biology**
 - hair and tissue mineral analysis
 - identification of trace metals in teeth
 - spectral fingerprinting of bacterial strains
 - identification of bacterial spores, molds, pollens and proteins
- **defense/homeland security**
 - detection of uranium in material,
 - high sensitivity detection of chemical and biological agents
 - *in situ* detection of land mines
- **forensic science**
 - identifying gunshot residue on hands
 - pen ink characterization
- **art conservation**
 - identifying pigments in paintings
 - dating/cleaning ancient marble



3 Current “Super-Stars” of Atomic Spectroscopy

1. electrothermal atomization-atomic absorption spectrometry (ETA-AAS)
2. inductively coupled plasma-atomic emission spectrometry (ICP-AES)
3. inductively coupled plasma-mass spectrometry (ICP-MS)



Suggested Reading

Comparing several atomic spectrometric methods to the super stars: special emphasis on laser induced breakdown spectrometry, LIBS, a future super star†

James D. Winefordner, Igor B. Gornushkin, Tiffany Correll, Emily Gibb, Benjamin W. Smith and Nicoló Omenetto

Department of Chemistry, University of Florida, Gainesville FL 32611, USA

J. Anal. At. Spectrom., 2004, 19, 1061–1083

This journal is © The Royal Society of Chemistry 2004



Review

Characterization of laser induced plasmas by optical emission spectroscopy: A review of experiments and methods

C. Aragón*, J.A. Aguilera

Departamento de Física, Universidad Pública de Navarra, Campus de Arrosadía, E-31006 Pamplona, Spain

Spectrochimica Acta Part B 63 (2008) 893–916



Advantages of LIBS

- 1) extremely fast analysis compared to competing technologies
- 2) multi-elemental analysis, light from all constituents collected without bias (see Periodic Table)
- 3) analysis can be performed at standoff distances
- 4) technique is applicable to all substrates (gas, solid, and liquid)
- 5) requires minimal or no sample prep
- 6) exquisite spatial resolution, $\sim 1 \mu\text{m}$
- 7) depth resolution (micron or sub-micron)



(■ Solids ■ Liquids ■ Gases ■ Artificially prepared)

| | | | | | | | | | | | | | | | | | |
|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|-----------------|----------------|-----------------|-----------------|
| 1 <u>H</u> | | | | | | | | | | | | | | | | | 2 <u>He</u> |
| 3 <u>Li</u> | 4 <u>Be</u> | | | | | | | | | | | 5 <u>B</u> | 6 <u>C</u> | 7 <u>N</u> | 8 <u>O</u> | 9 <u>F</u> | 10 <u>Ne</u> |
| 11 <u>Na</u> | 12 <u>Mg</u> | | | | | | | | | | | 13 <u>Al</u> | 14 <u>Si</u> | 15 <u>P</u> | 16 <u>S</u> | 17 <u>Cl</u> | 18 <u>Ar</u> |
| 19 <u>K</u> | 20 <u>Ca</u> | 21 <u>Sc</u> | 22 <u>Ti</u> | 23 <u>V</u> | 24 <u>Cr</u> | 25 <u>Mn</u> | 26 <u>Fe</u> | 27 <u>Co</u> | 28 <u>Ni</u> | 29 <u>Cu</u> | 30 <u>Zn</u> | 31 <u>Ga</u> | 32 Ge | 33 <u>As</u> | 34 Se | 35 <u>Br</u> | 36 <u>Kr</u> |
| 37 <u>Rb</u> | 38 <u>Sr</u> | 39 <u>Y</u> | 40 <u>Zr</u> | 41 Nb | 42 <u>Mo</u> | 43 Tc | 44 <u>Ru</u> | 45 <u>Rh</u> | 46 <u>Pd</u> | 47 <u>Ag</u> | 48 <u>Cd</u> | 49 <u>In</u> | 50 <u>Sn</u> | 51 <u>Sb</u> | 52 Te | 53 <u>I</u> | 54 <u>Xe</u> |
| 55 <u>Cs</u> | 56 <u>Ba</u> | 57 <u>La</u> | 72 <u>Hf</u> | 73 <u>Ta</u> | 74 <u>W</u> | 75 <u>Re</u> | 76 Os | 77 <u>Ir</u> | 78 <u>Pt</u> | 79 <u>Au</u> | 80 <u>Hg</u> | 81 <u>Tl</u> | 82 <u>Pb</u> | 83 Bi | 84 Po | 85 At | 86 <u>Rn</u> |
| 87 Fr | 88 Ra | 89 Ac | 104 Rf | 105 Db | 106 Sg | 107 Bh | 108 Hs | 109 Mt | 110 Uun | 111 Uuu | 112 Uub | | 114 Uuq | | 116 Uuh | | |
| | | | 58 <u>Ce</u> | 59 Pr | 60 <u>Nd</u> | 61 Pm | 62 <u>Sm</u> | 63 <u>Eu</u> | 64 <u>Gd</u> | 65 Tb | 66 Dy | 67 Ho | 68 <u>Er</u> | 69 Tm | 70 Yb | 71 Lu | |
| | | | 90 <u>Th</u> | 91 Pa | 92 <u>U</u> | 93 Np | 94 <u>Pu</u> | 95 Am | 96 Cm | 97 Bk | 98 Cf | 99 Es | 100 Fm | 101 Md | 102 No | 103 Lr | |

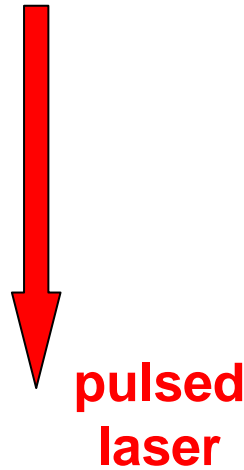
www.arl.army.mil/wmrd/LIBS

The LIBS Process

1. laser interaction with the target
2. removal of samples mass (ablation)
3. plasma formation (breakdown)
4. element specific emission

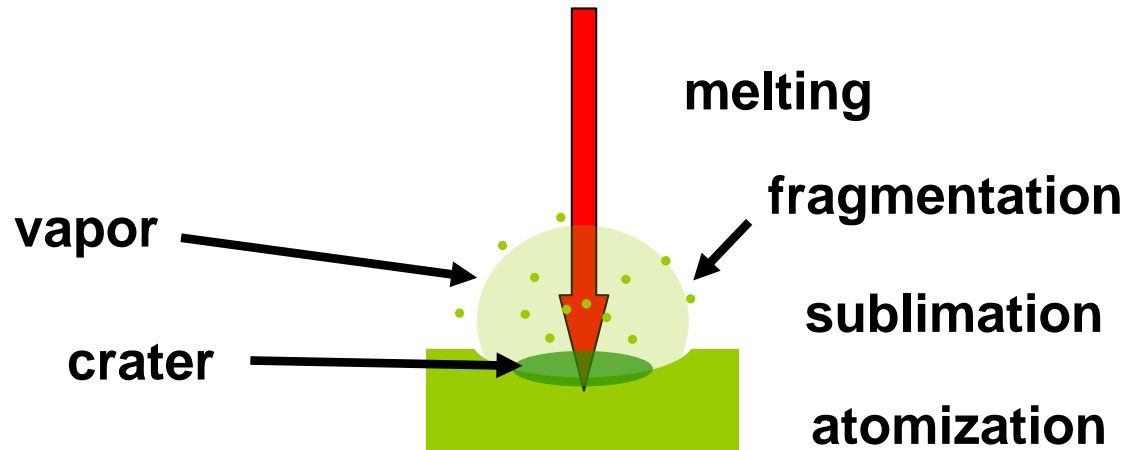


1) laser interaction with the target



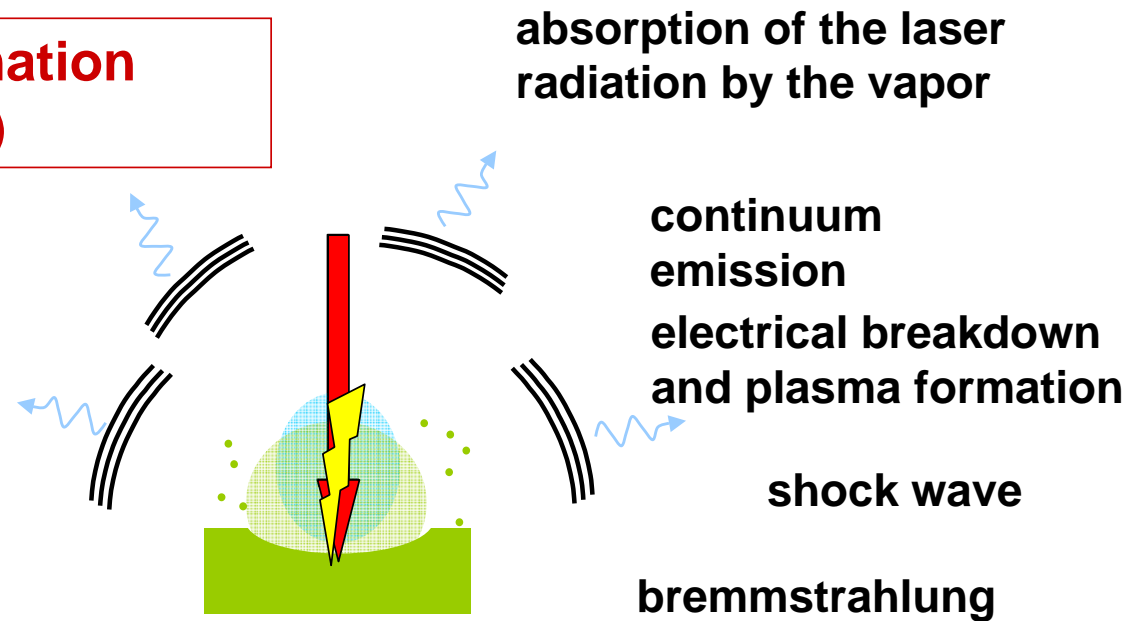
- initiated by absorption of energy by the target from a pulsed radiation field.
- pulse durations are on the order of nanoseconds, but LIBS has been performed with pico- and femto-second laser pulses.

2) removal of samples mass (ablation)



- absorbed energy is rapidly converted into heating, resulting in vaporization of the sample (ablation) when the temperature reaches the boiling point of the material.
- removal of particulate matter from the surface leads to the formation of a vapor above the surface.

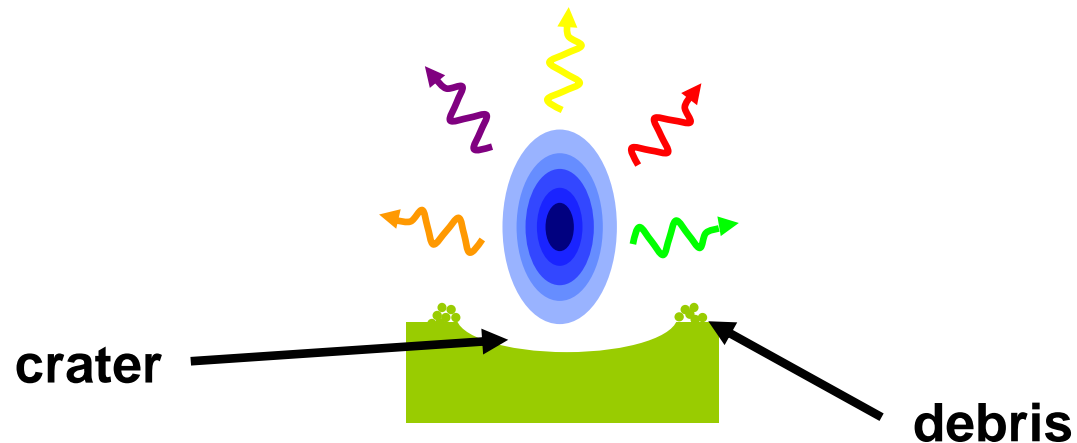
3) plasma formation (breakdown)



- The laser pulse continues to illuminate the vapor plume.
- The vapor condenses into sub-micrometer droplets that lead to absorption and scattering of the laser beam, inducing strong heating, ionization, and plasma formation.

4) expansion and element specific emission (atomic or ionic)

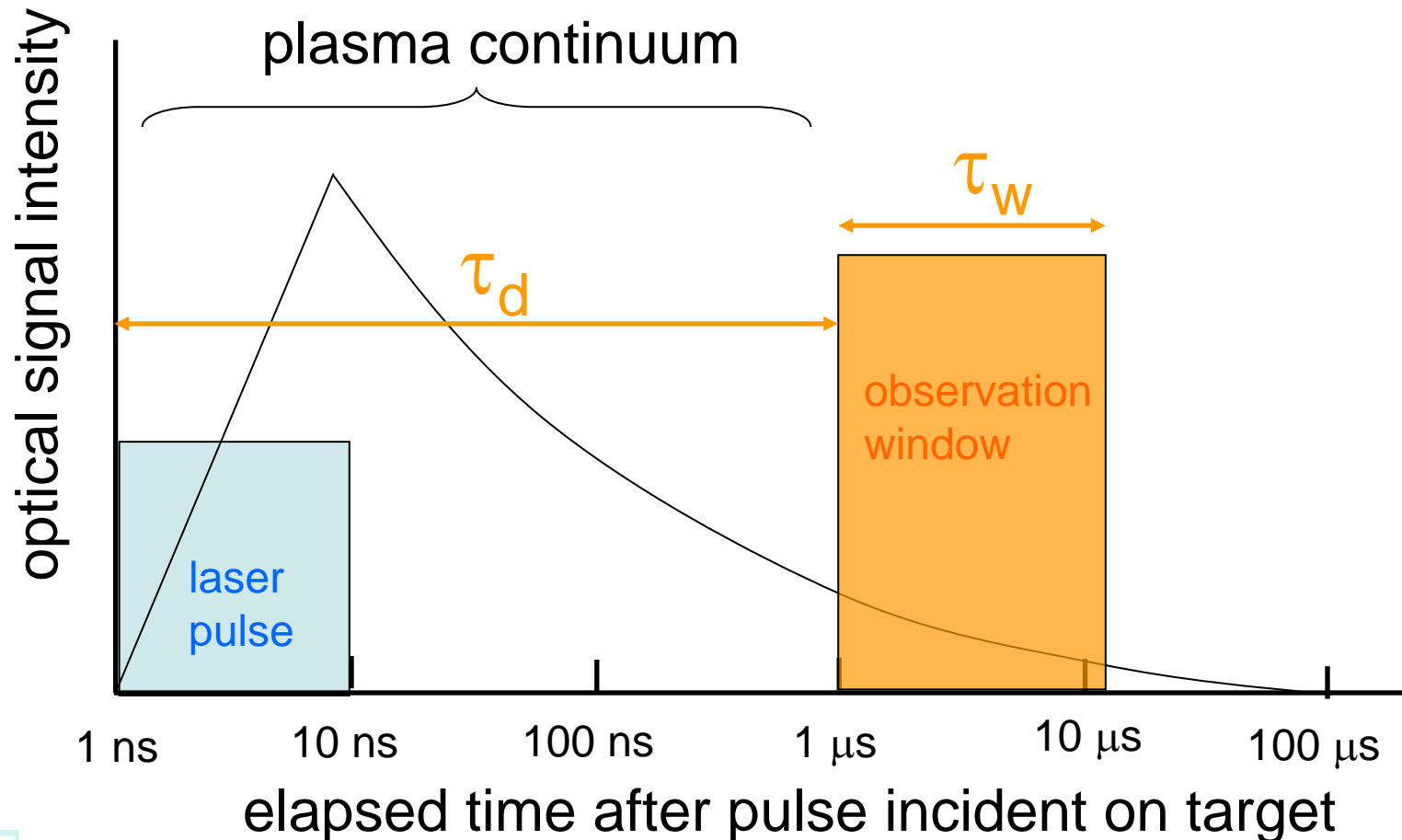
spontaneous emission
as atoms/ions decay to
ground state



- The dynamical evolution of the plasma plume is then characterized by a fast expansion and subsequent cooling.
- Approximately 1 microsecond after the ablation pulse, spectroscopically narrow atomic/ionic emissions may be identified in the spectrum.

Temporal History of a LIBS Plasma

D.A. Cremers and L.J. Radziemski, in "Handbook of Laser-Induced Breakdown Spectroscopy," (2006) p. 24



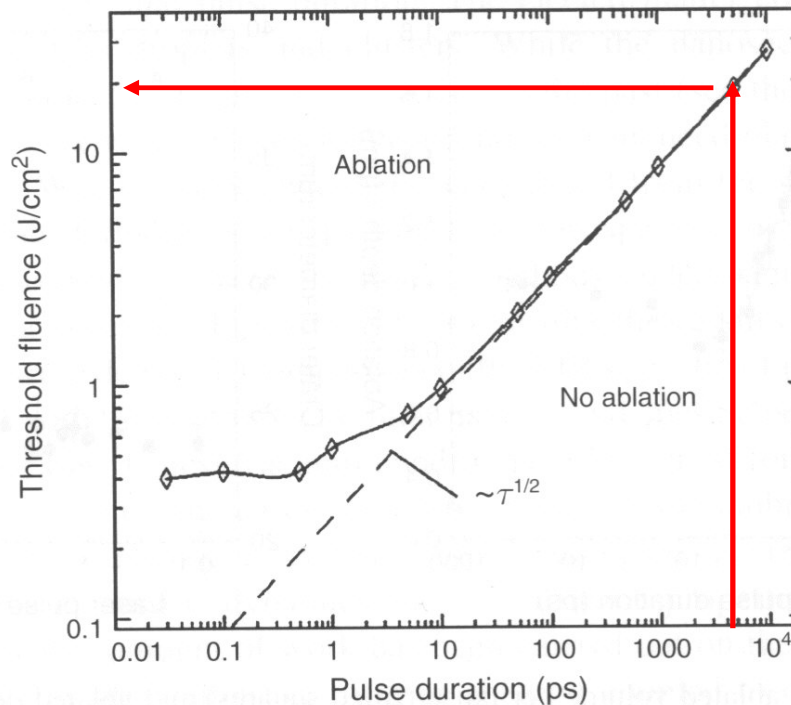
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1. laser interaction with the target

M. Sabsabi, in "Laser-Induced Breakdown Spectroscopy," (2007) p. 159



- plot for aluminum
- typical ns lasers
 - assume a 100 micron spot size (modest focusing)
 - requires approximately 1 mJ (modest pulse energy)
 - to be well above ablation threshold, typically 10's of mJ used

Fig. 5. Threshold fluence for ablation as a function of the pulse duration.

1. laser pulse fluence

- so a 1 mJ ns pulse focused to a spot size of 100 μm yields a **fluence** of $\sim 10 \text{ J/cm}^2$
- assuming a 10 ns pulse (typical) operating at 10 Hz (typical)

$$10 \frac{\text{J}}{\text{cm}^2} \cdot \frac{1 \text{ pulse}}{10 \text{ ns}} \cdot 10 \text{ pulses} = 10 \frac{\text{GW}}{\text{cm}^2}$$

- so **irradiances** of 1-10 GW/cm^2 are typical



1. laser irradiance calculation

$$I_{\min} = \frac{\rho L_V \kappa^{1/2}}{\Delta t^{1/2}} \text{ (W/cm}^2\text{)}$$

- ρ = density
- L_V = latent heat of vaporization
- κ = thermal diffusivity
- Δt = laser pulse length

- $I_{\min} A_1 = 1.75 \times 10^8 \text{ W/cm}^2$
 - for a 10 ns pulse, focused to a 100 μm spot: $\sim 130 \mu\text{J}$



1. laser pulse wavelength

- all wavelengths have been used / can be used
- different authors have reported different dependence of breakdown threshold fluence with wavelength
 - (some increasing, some decreasing)
 - [see *E. Tognoni et al. Spectrochimica Acta B* **57** (2002) 1115-1130]



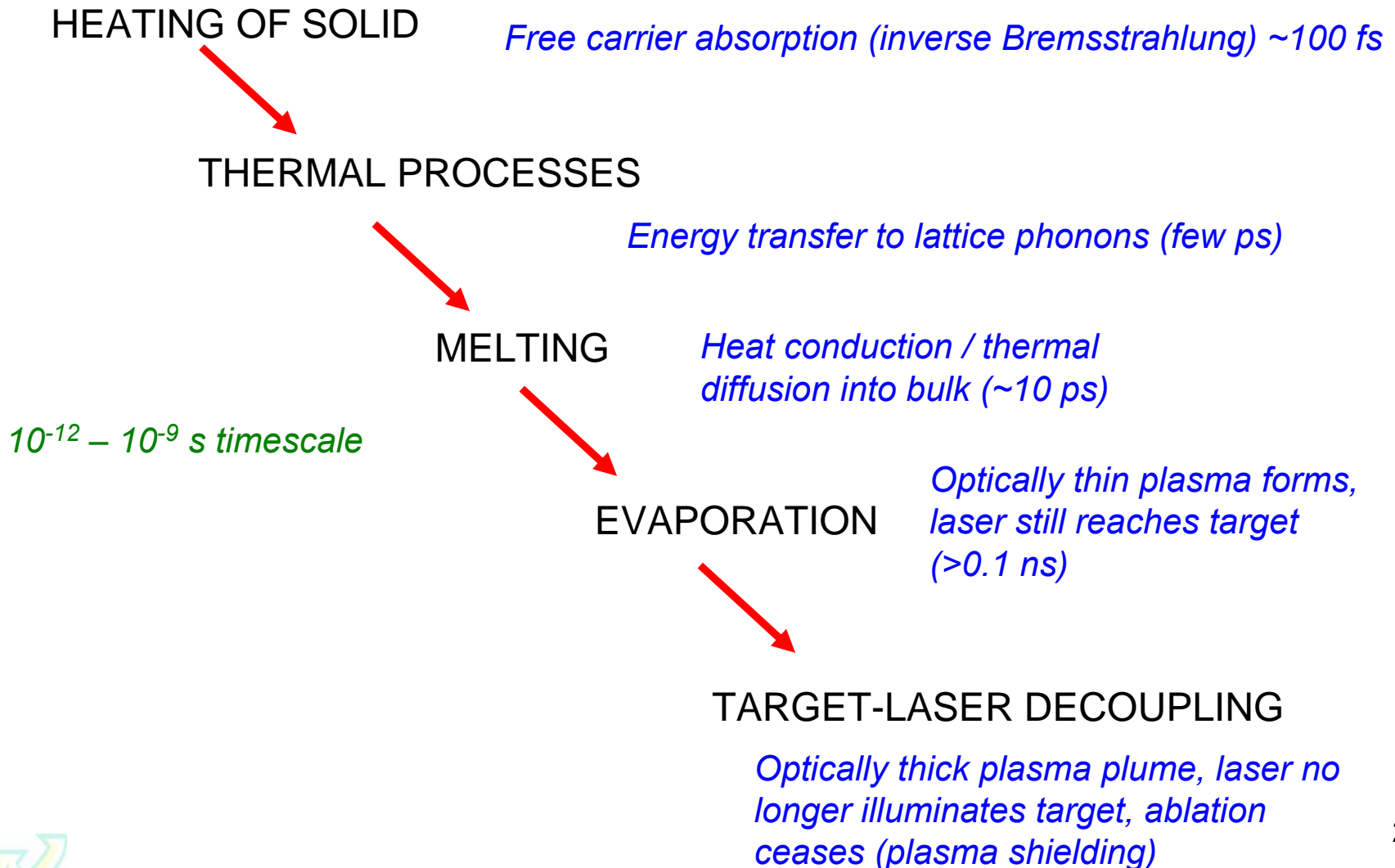
1. laser pulse wavelength

- complexity arises from:
 - the multiple steps required to form a laser-induced plasma
 - the variety of target properties (reflectivity, dielectric constant)
 - what do you want out of it?
 - most mass removed
 - highest intensity
 - best SNR
- “In addition, wavelength is an important factor for the initiation of plasma. Short laser wavelengths aid the ablation process because a good absorption of laser irradiance by the material occurs. Nevertheless, the inverse Bremsstrahlung process is more efficient with IR wavelength than with UV radiation and the laser induced breakdown in ambient air is then favored at larger wavelengths.”
L.M. Cabalin and J.J. Laserna, Spectrochimica Acta Part B 53 (1998) 723-730



2. removal of samples mass (ablation)

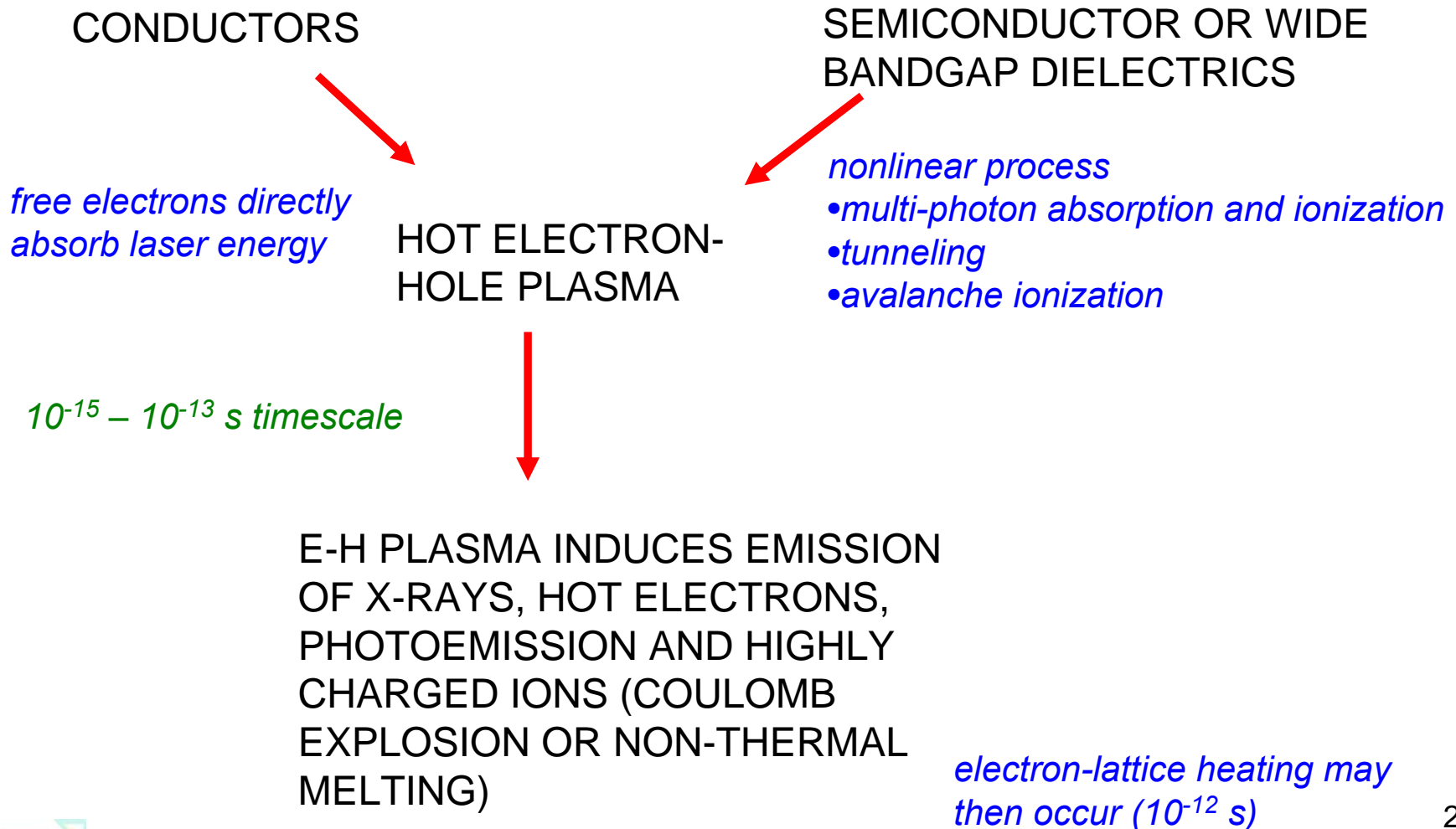
for solid targets and ns pulses...



2. removal of samples mass (ablation)

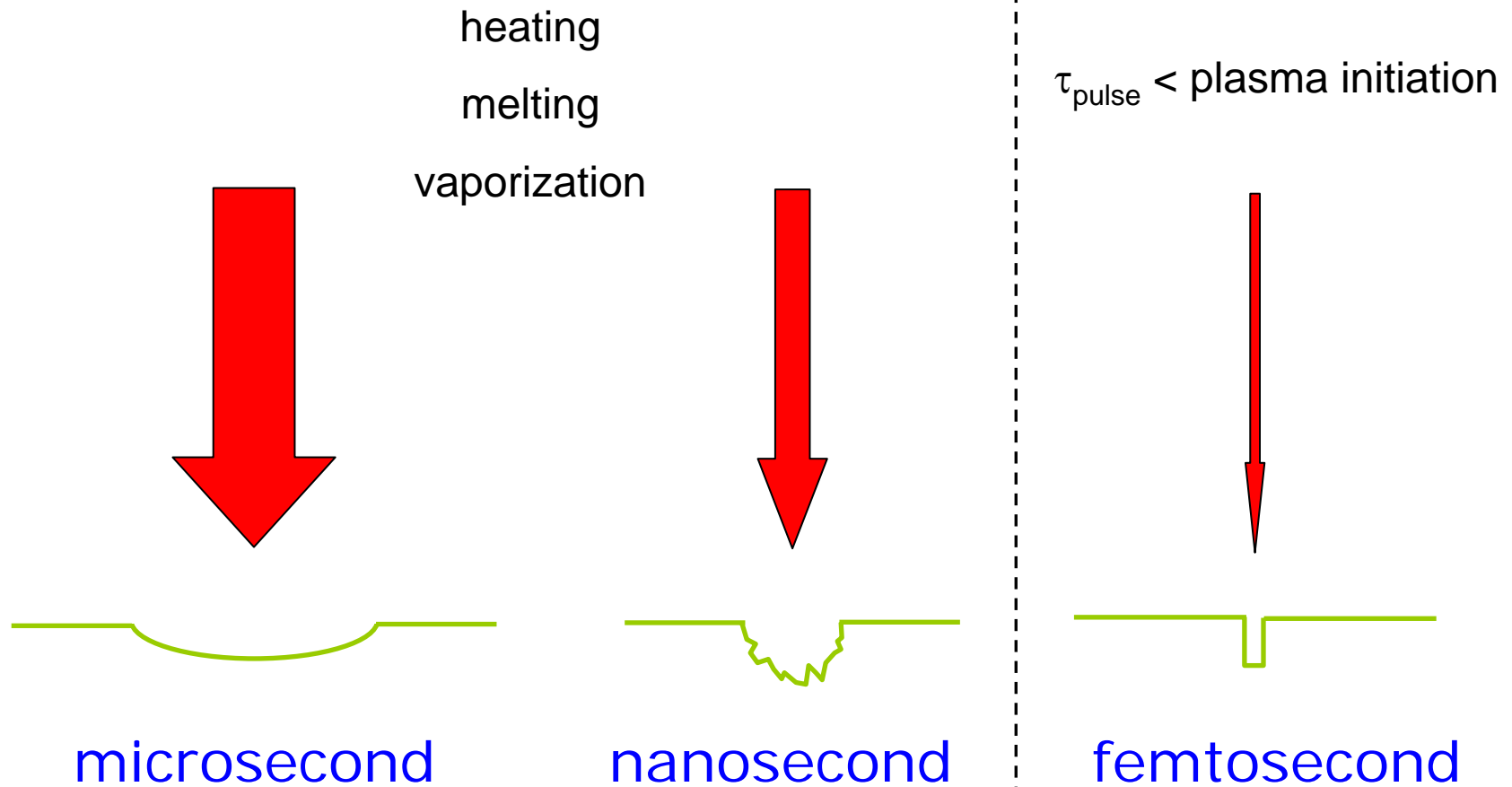
for solid targets and fs pulses...

R.E. Russo, in "Laser-Induced Breakdown Spectroscopy," (2007) p. 49



Physics of Plasma Formation: ablation

D.A. Cremers and L.J. Radziemski, in "Handbook of Laser-Induced Breakdown Spectroscopy," (2006) p. 45



2. removal of samples mass (ablation)

B.N. Chichkov et al., Appl. Phys. A63, 109-115 (1996)

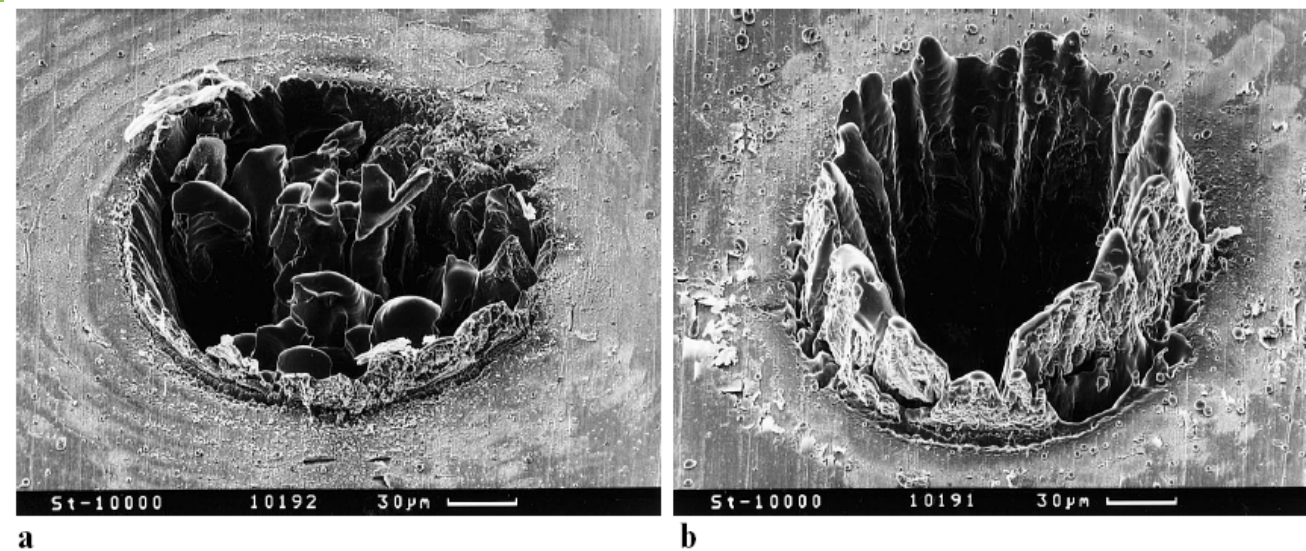


Fig. 2a, b. Schematic of nanosecond-pulse laser ablation and holes drilled in a 100 μm thick steel foil with (a) 80 ps, 900 μJ , $F = 3.7 \text{ J/cm}^2$; and (b) 3.3 ns, 1 mJ, and $F = 4.2 \text{ J/cm}^2$ laser pulses at 780 nm

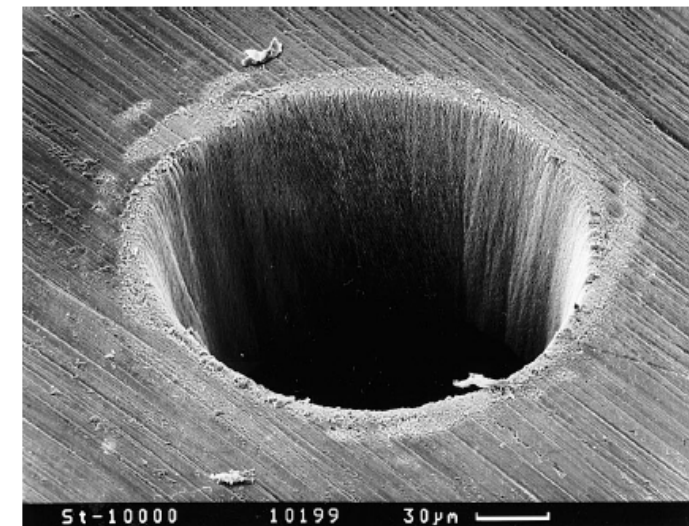


Fig. 1. Schematic of femtosecond-pulse laser ablation and a SEM photograph of a hole drilled in a 100 μm thick steel foil with 200 fs, 120 μJ , $F = 0.5 \text{ J/cm}^2$ laser pulses at 780 nm

ablation efficiency

e.g., A. Semerok et al., *Appl. Surf. Sci.* **311**, 138-139 (1999).

B. Salle et al., *Appl. Phys. A* **69**, S381-383 (1999).

$$\text{Definition 1} = \frac{\text{volume of matter ablated}}{\text{laser pulse energy}}$$

$$\text{Definition 2} = \frac{\text{crater depth}}{\text{laser fluence}}$$

If the laser intensity distribution corresponds to the crater profile, the two definitions are equivalent.



typically ng to μg of material ablated

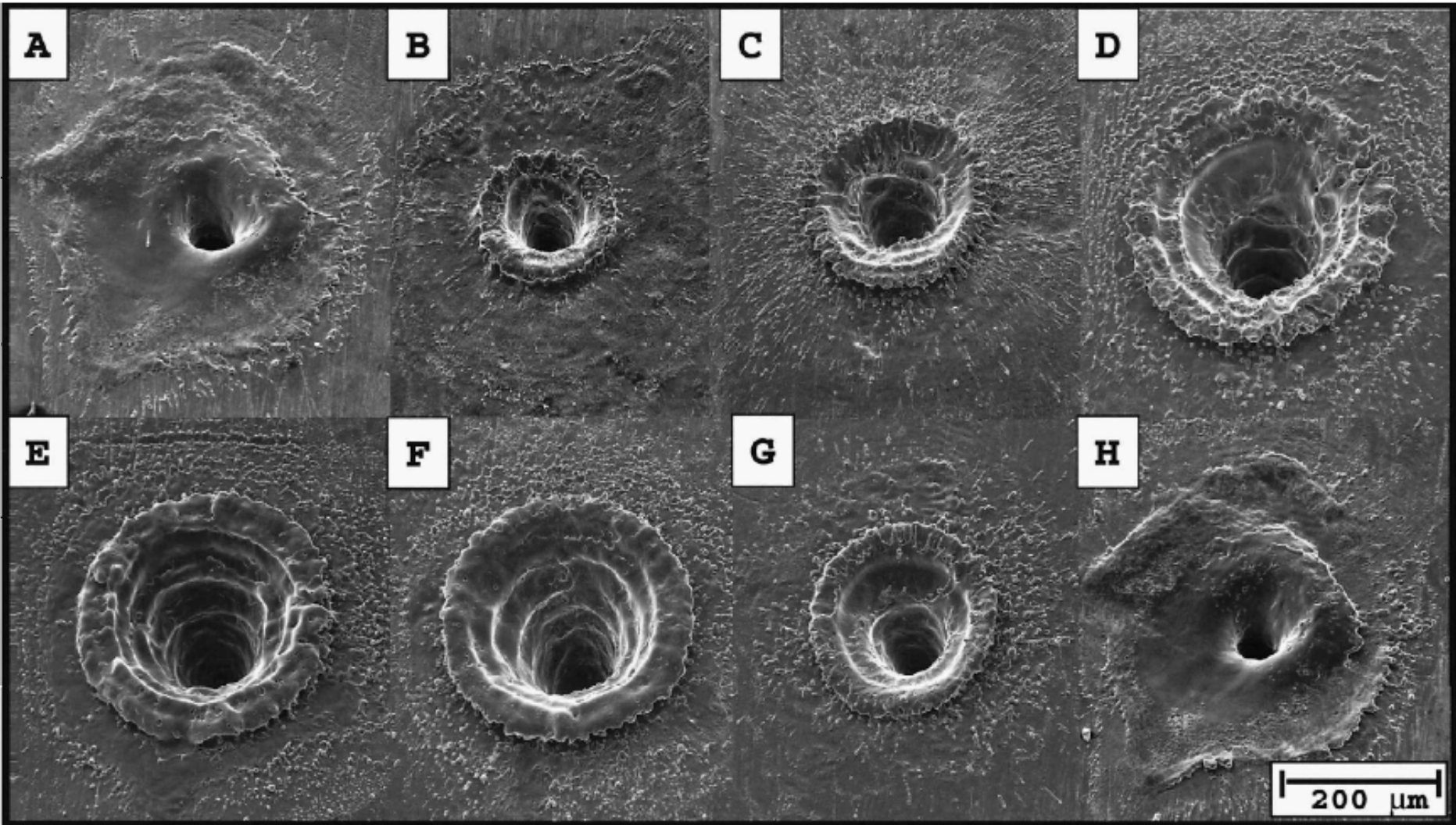
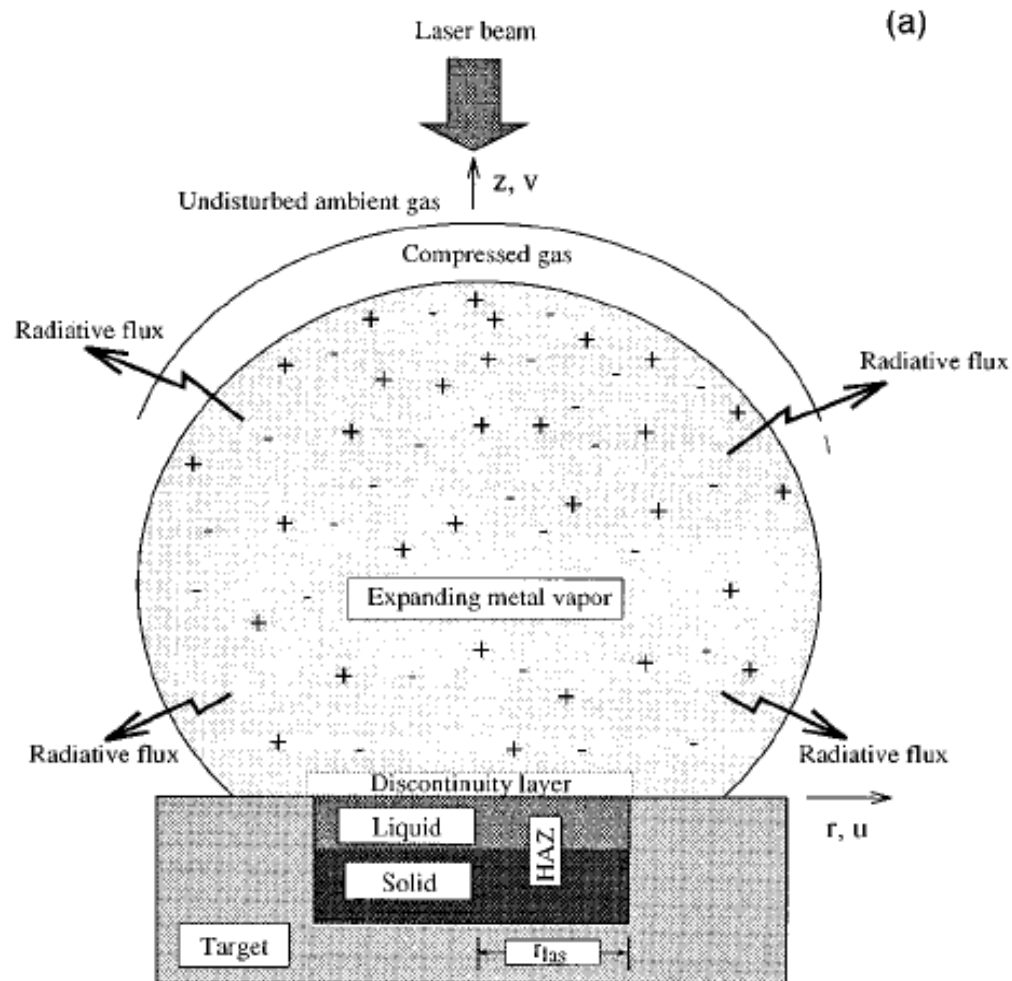


FIG. 4. SEM images of the ablation craters vs. laser pulse delay for a steel sample using (A) single-pulse, and dual-pulse delays of (B) 0.3 μs , (C) 0.5 μs , (D) 2.5 μs , (E) 25 μs , (F) 50 μs , (G) 100 μs , and (H) 300 μs .

3. plasma formation (breakdown)

J.R. Ho et al., J. Appl. Phys. 79, 7205-7215 (1996)



3. plasma formation (breakdown)

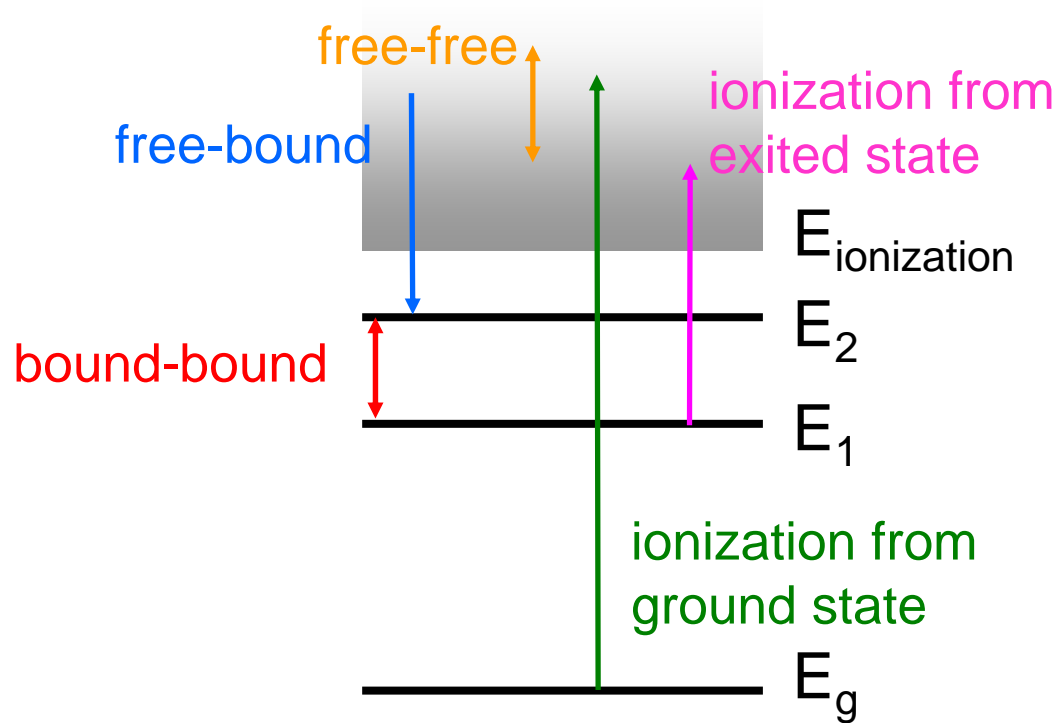
Problem: how do photons of relatively low energy, 1-2 eV, (compared to ionization threshold of common gases) generate a breakdown (a plasma at 10,000-50,000K) ?

Three distinct but overlapping stages:

1. plasma ignition
2. plasma growth (electron avalanche or cascade) and interaction with laser pulse
3. plasma development accompanied by shock wave generation and propagation (“breakdown”)



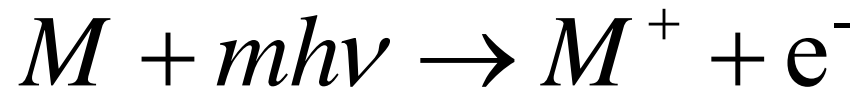
Transitions in an Atom or Ion



3. plasma formation (breakdown)

1. cascade or avalanche requires an initial electron

- multiphoton absorption/ionization

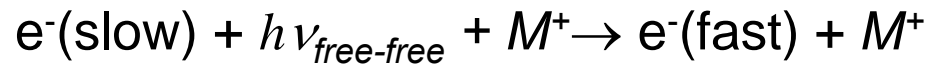


- local radioactivity
- cosmic rays
- local free electron density



3. plasma formation (breakdown)

2. electron cascade or avalanche occurs by inverse bremsstrahlung (free-free absorption)

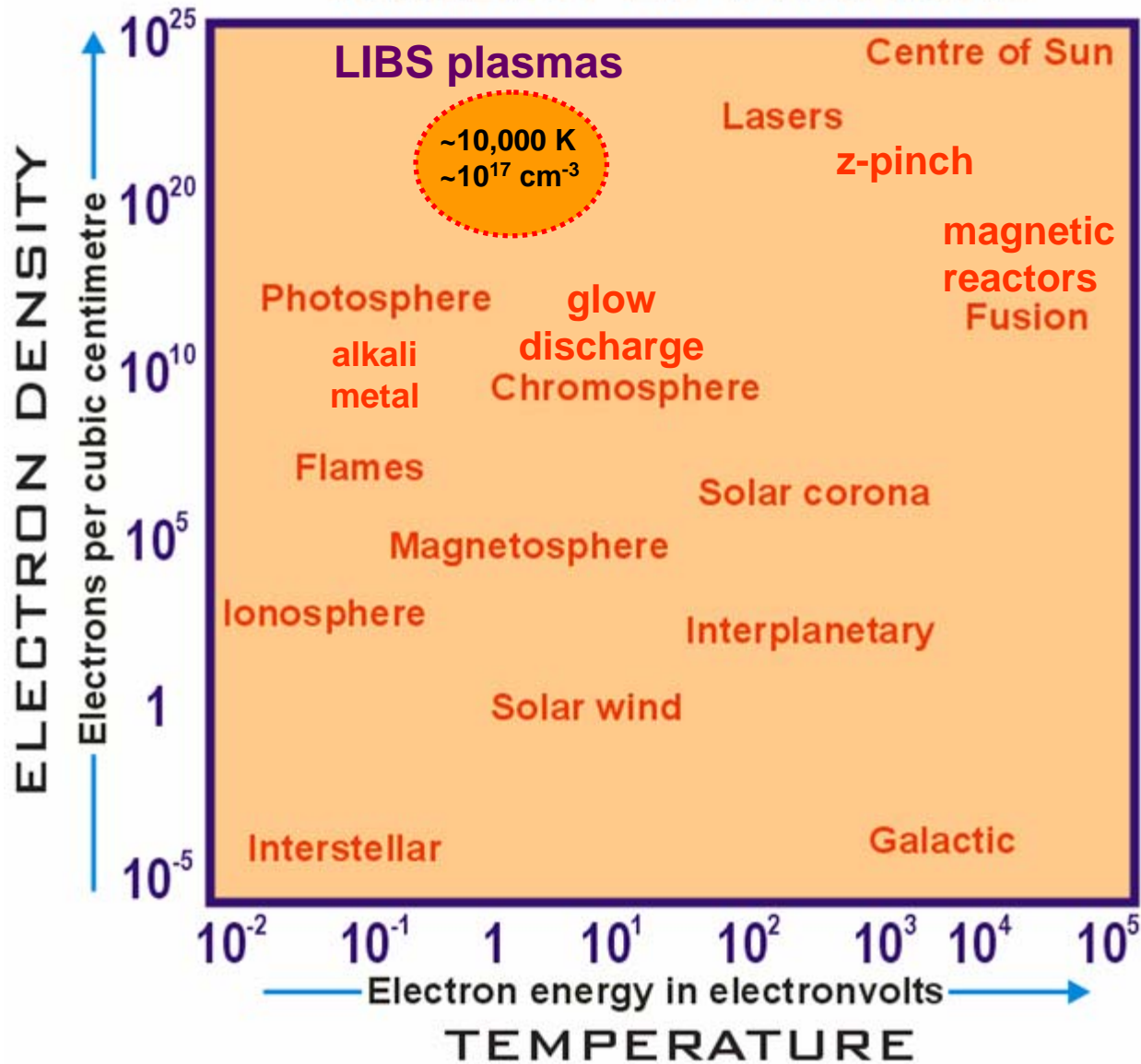


- electrons absorb photons from laser field (in the presence of gas) for momentum transfer between collisions with neutral species
- acquire sufficient energy for collisional ionization of gas atoms
- electron density increases exponentially via cascade

$$n_e \sim 1-10 \text{ cm}^{-3} \rightarrow 10^{17}-10^{20} \text{ cm}^{-3}$$



RANGES OF PLASMAS



3. plasma formation (breakdown)

3. “breakdown” is arbitrarily defined

$$n_e \sim 10^{13} \text{ cm}^{-3} \text{ or degree of ionization of } 10^{-3}$$

permits significant absorption and scattering of incident laser beam leads very fast to a fully developed plasma and shockwave

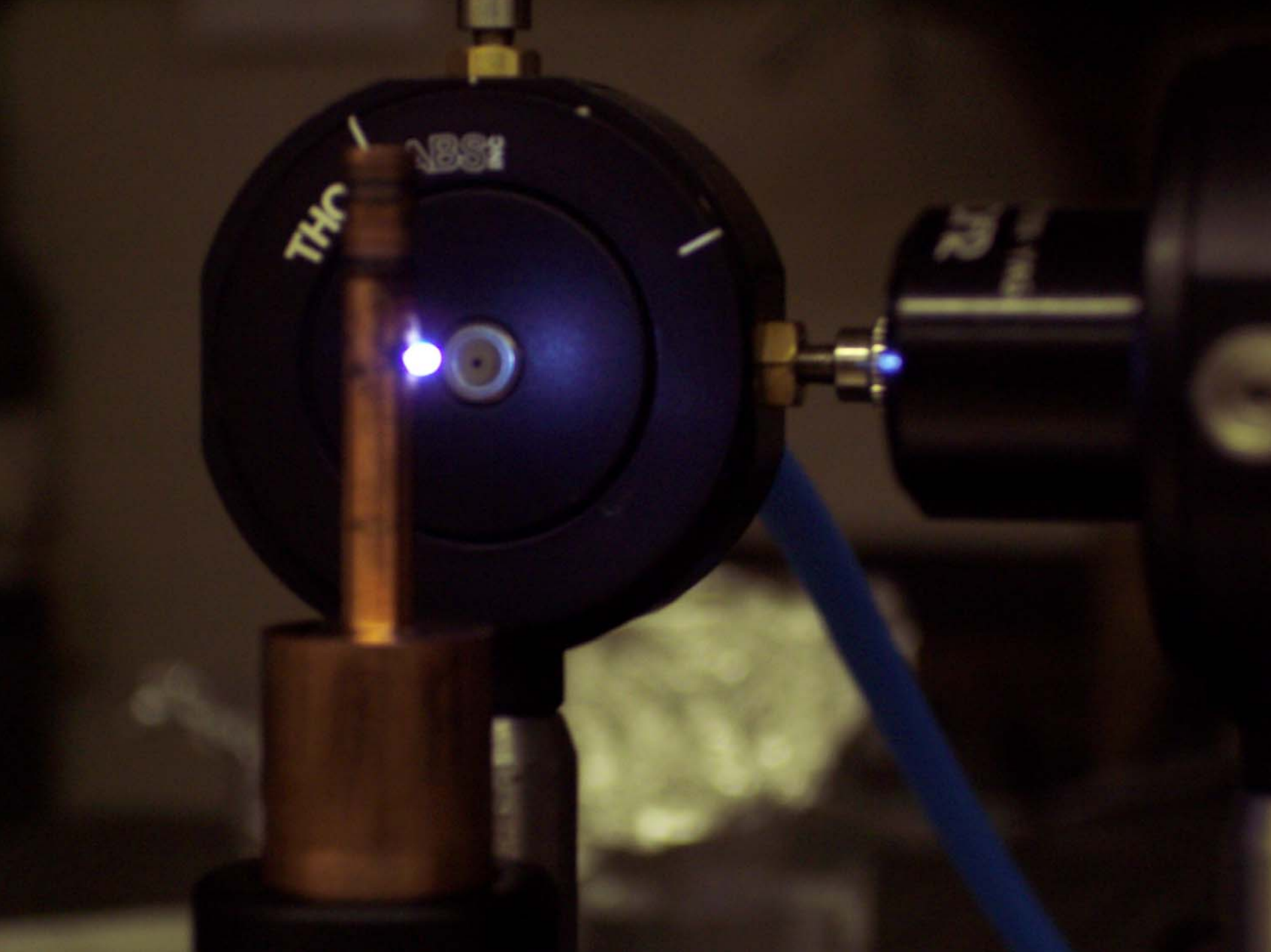
$$10^{13} \text{ cm}^{-3} \rightarrow 10^{17} - 10^{20} \text{ cm}^{-3}$$





**A Stark look at
plasma breakdown**





3. plasma formation (plasma shielding)

eventually, the plasma becomes opaque to the laser beam and the target is shielded

occurs when plasma frequency becomes greater than the laser frequency

$$\omega_p \approx \omega$$

or when

$$n_e \sim \left(10^{21} / \lambda^2\right) \text{cm}^{-3}$$

λ in micron



4. Expansion and Emission

J.F. Ready, "Effect of High Power Laser Radiation," 1971

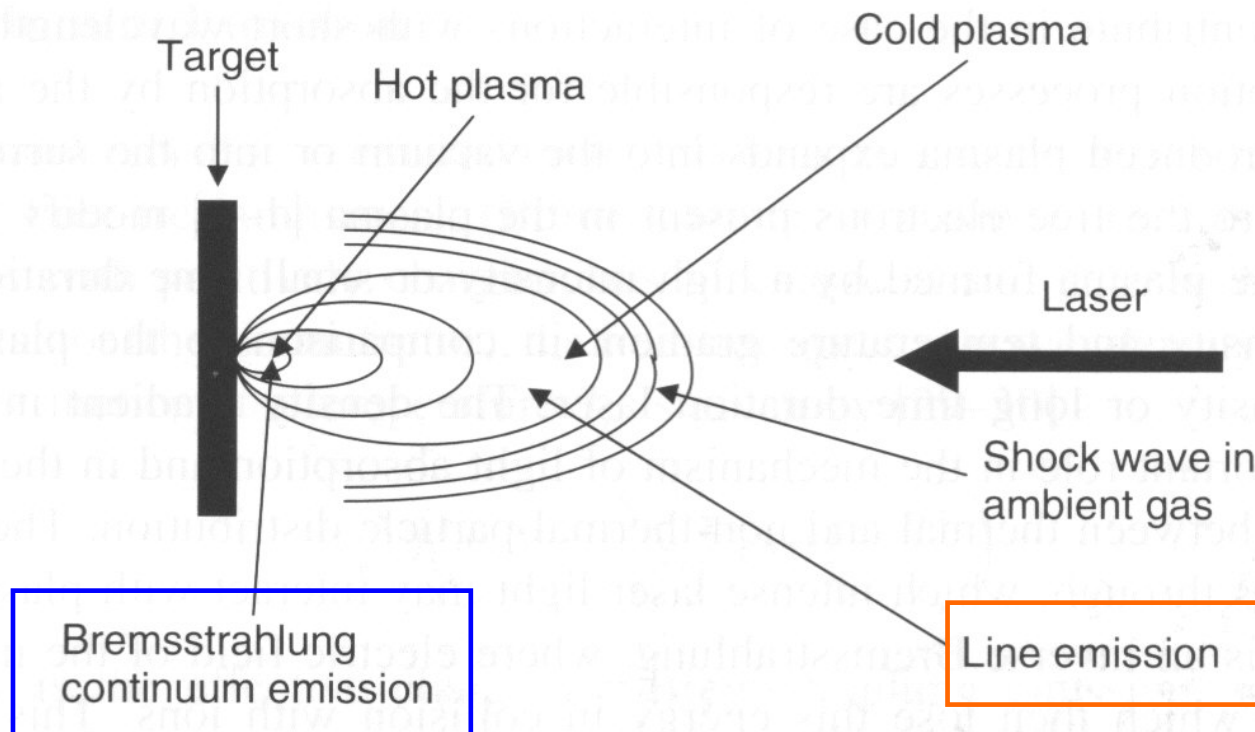
- laser absorption in the expanding vapor/plasma can generate three different types of waves depending on irradiation intensity
 - laser-supported combustion (LSC)
 - laser-supported detonation (LSD)
 - laser-supported radiation (LSR)



expanding/cooling plasma

V.N. Rai and S.N. Thakur, in "Laser-Induced Breakdown Spectroscopy,"(2007) p. 85

- the dynamic nature of this emission source is why time gating is used



The Goal of LIBS Plasma Creation

- to create an optically thin plasma which is in thermodynamic equilibrium and whose elemental composition is the same as that of the sample
 - if achieved, spectral line intensities can be connected to relative concentrations of elements
 - typically these conditions are only met *approximately*.



ESAWIN - [Spectrum Spe-Viewer - 062706C_10]

File View Configuration Measurement ROI Service File Service Window Help

On-Chip Accus 10 Gate Width 1000 ns Gate Delay 1500 ns Amplific. 2500 Apply

Set Global

Activate Release Spectra Release Set

Root Name file Number 1

25832
20666
15499
10333
5166
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200.000 264.200 328.400 392.600 456.800 521.000 585.200 649.400 713.600 777.800 842.000

712.314 nm I: 26.603 X: 164.6 Y: 892.6 Ord: 34 Element: Nb

Ready No Measure Expo: 1000 ms

456.423

start Ready No Measure Expo: 1000 ms

start Ready No Measure Expo: 1000 ms

start Ready No Measure Expo: 1000 ms



Spectral Line Radiant Intensity

$$I = \frac{h\nu gAN}{4\pi} = \left(\frac{hcN_0 gA}{4\pi\lambda Z} \right) \exp\left(-\frac{E}{kT}\right)$$

I = intensity (given in units of W/sr)

g = statistical weight of level

A = Einstein A coefficient

N_0 = total species population

Z = partition function (statistical weight of ground state)

E = Energy of upper state of transition

Low concentrations required

J.A. Aguilera et al., Spectrochim. Acta B 58, 221-237 (2003)

“Curve of growth”:

- for weak optically thin lines, $I \propto N$
- as abundance increases the plasma can become optically thick and the curve saturates

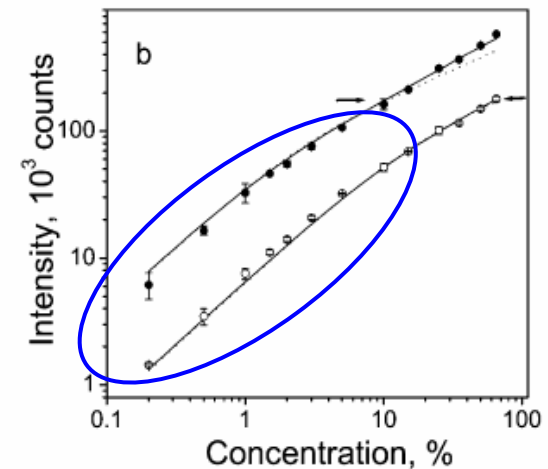
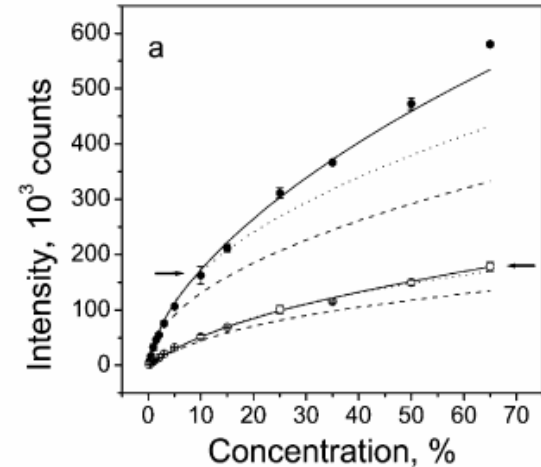
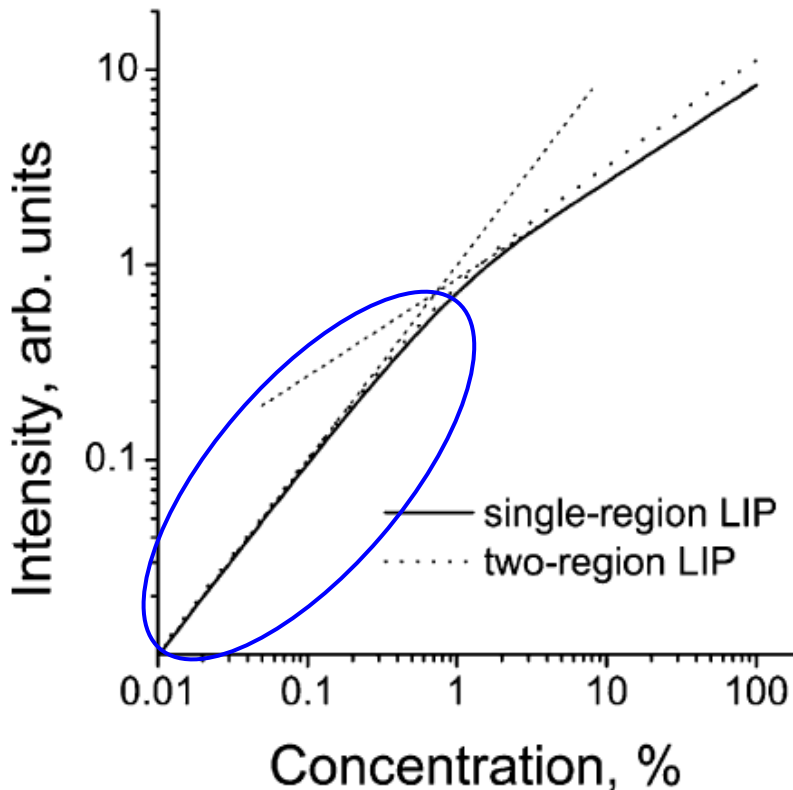


Fig. 4. Experimental and theoretical curves of growth in linear (a) and double logarithmic (b) plots of the Fe I spectral lines A (solid circles) and C (open circles). The theoretical curves have been calculated using the two-region model (solid lines)

Plasma Diagnostics: Temperature

$$I = \frac{h\nu gAN}{4\pi} = \left(\frac{hcN_0 gA}{4\pi\lambda Z} \right) \exp\left(-\frac{E}{kT}\right)$$

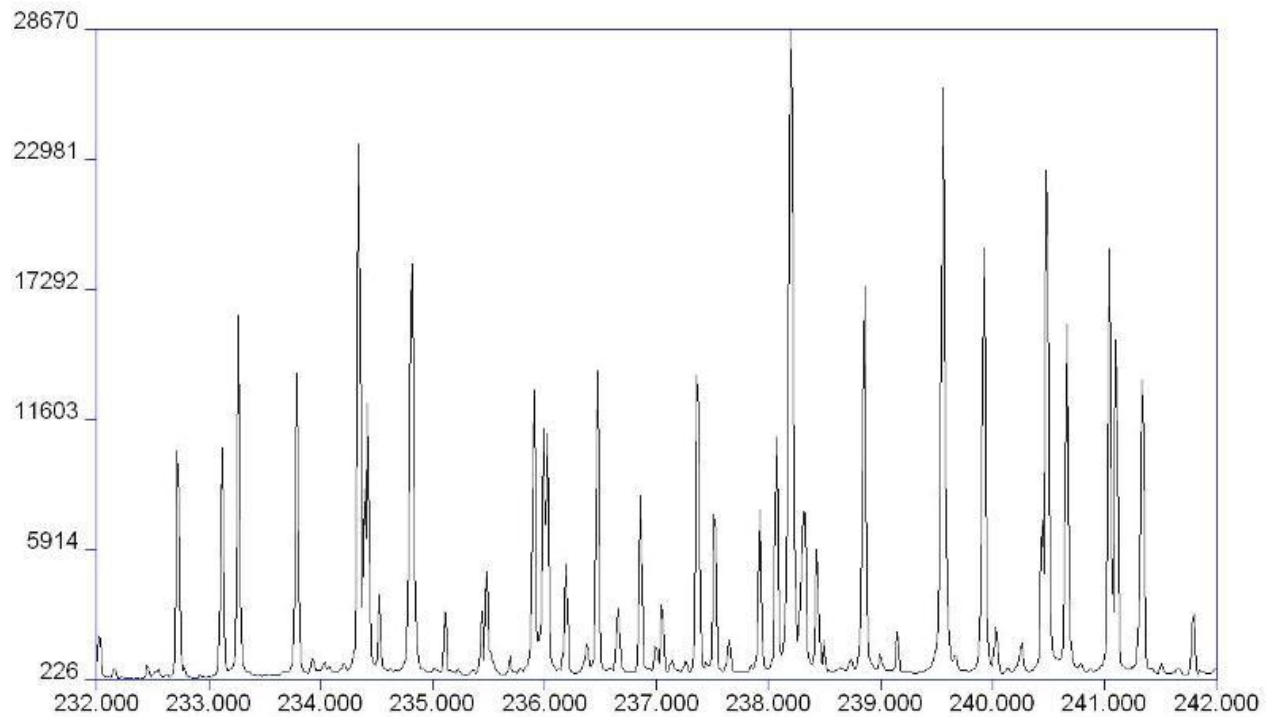
- better to write...

$$\ln\left(\frac{I\lambda}{gA}\right) = -\frac{E}{kT} - \ln\left(\frac{4\pi Z}{hcN_0}\right)$$

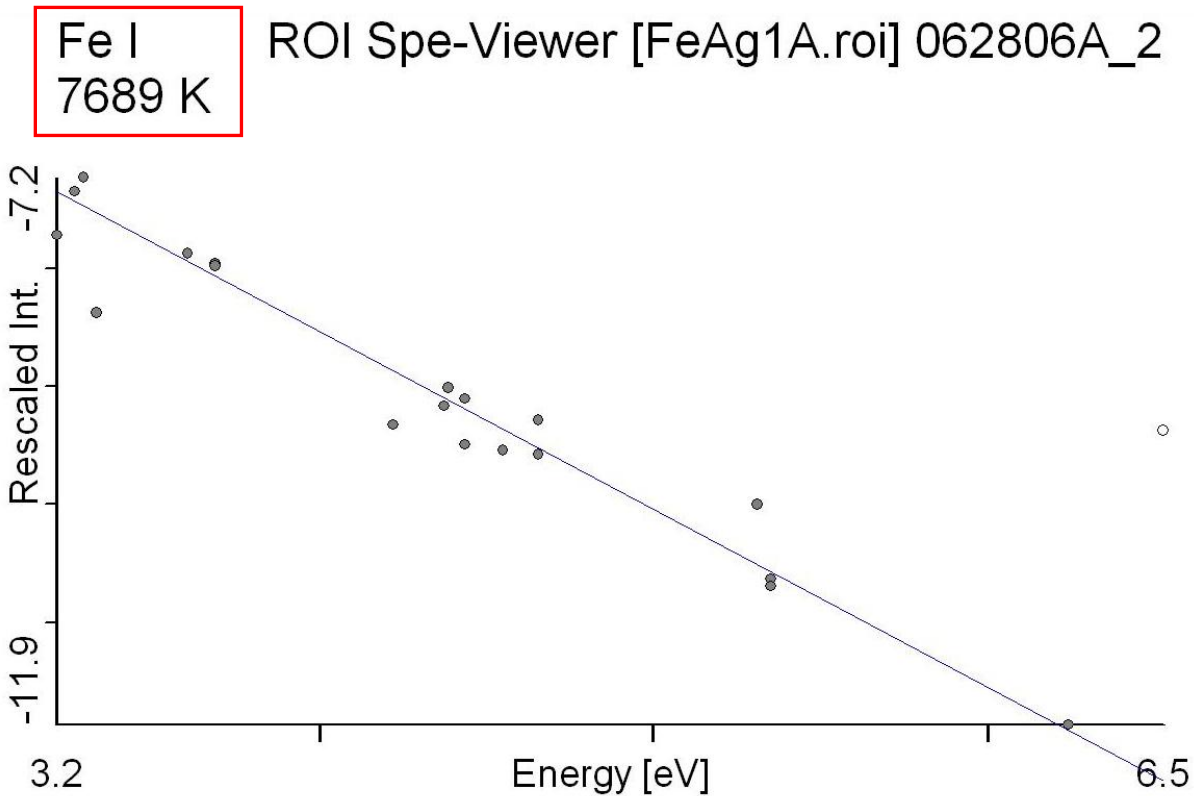
- This is a straight line with slope of $-1/kT$!
- So if we plot the adjusted measured line intensity vs. the upper state energy of transitions we can measure T of our plasma.



Fe₂O₃ / Ag Mixture



Fe Temperature



Boltzmann plot for 22 Fe transitions

Plasma Diagnostics

Temperature

plasma on water surface

Temperatures

calculated from H_β / H_γ

intensity ratio using

Boltzmann equation:

$$\frac{I_1}{I_2} = \frac{g_1 A_1}{g_2 A_2} \frac{\lambda_2}{\lambda_1} \exp\left(-\frac{|E_1 - E_2|}{kT_e}\right)$$



Plasma Diagnostics

electron density

FWHM of Stark-broadened lines used to calculate electron density N_e

$$N_e = C(N_e, T) \Delta\lambda_{FWHM}^{3/2}$$

Local Thermodynamic Equilibrium (LTE) can be assumed if

$$N_e \left(\text{cm}^{-3} \right) \geq 1.6 \times 10^{12} (\Delta E)^3 (T)^{\frac{1}{2}}$$

ΔE = energy difference with ground level (eV)

T = plasma temperature (K)



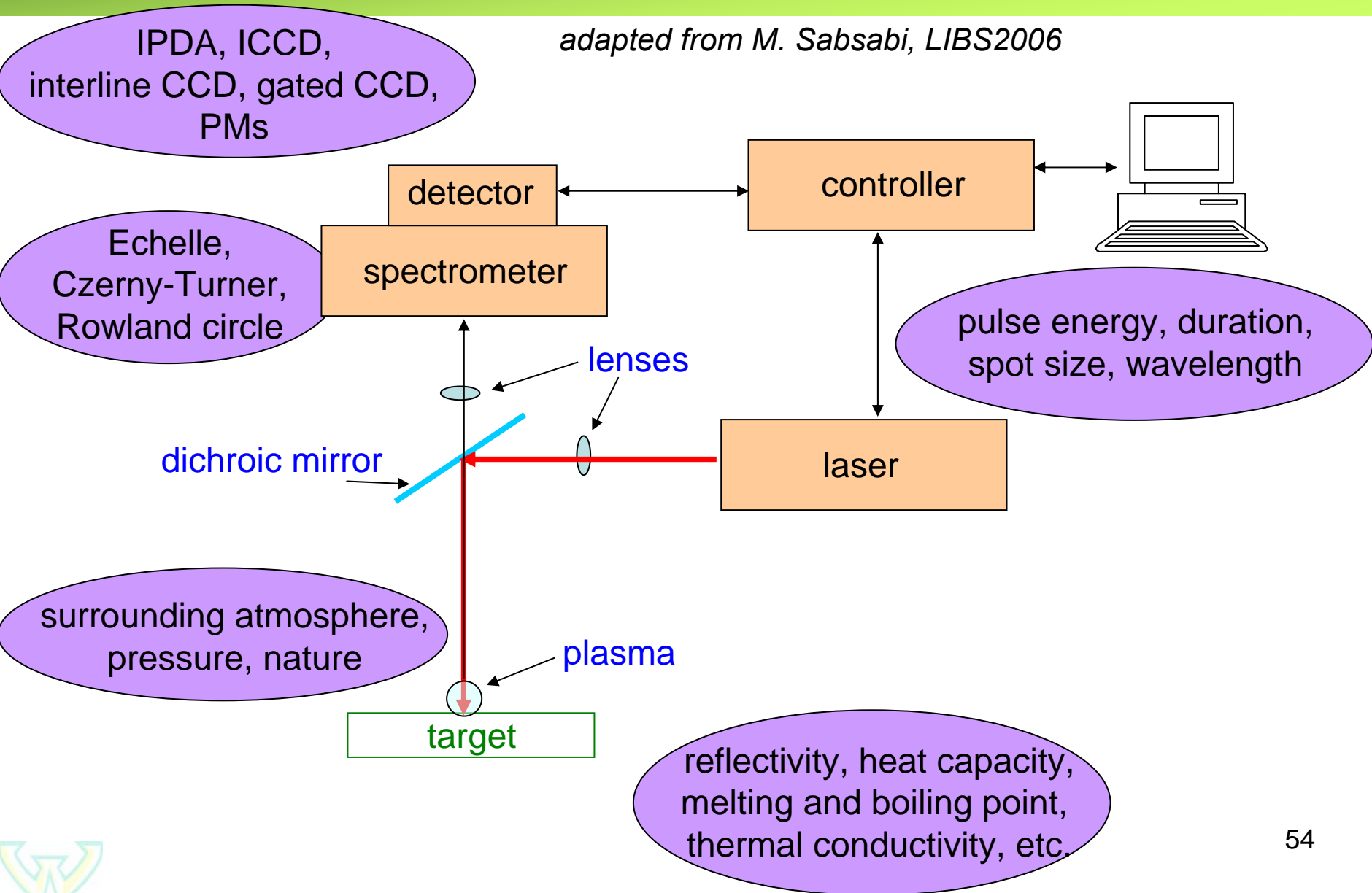
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LIBS components

adapted from M. Sabsabi, LIBS2006



Microchip lasers

K. Amponsah-Manager et al., Microchip laser ablation of metals: investigation of the ablation process in view of its application to laser-induced breakdown spectroscopy, JAAS 20, 544-551 (2005)

A. Freedman et al., Aluminum alloy analysis using microchip-laser induced breakdown spectroscopy, Spectrochim. Acta B 60, 1076-1082 (2005)

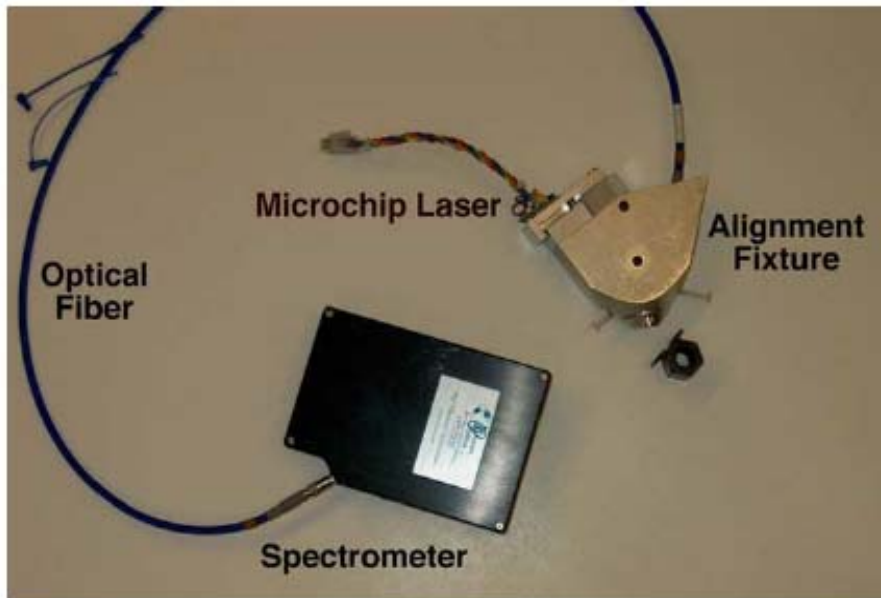


Fig. 1. Photograph of LIBS apparatus.

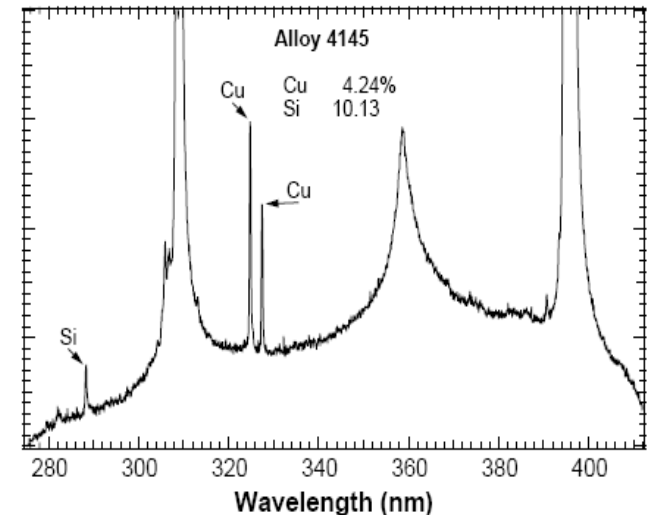
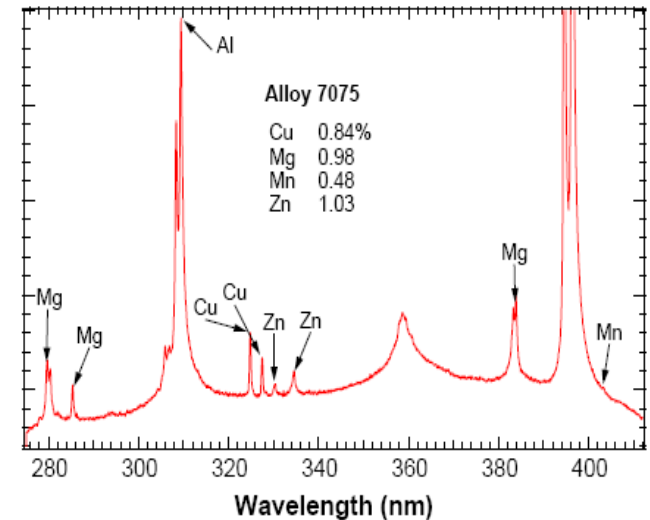
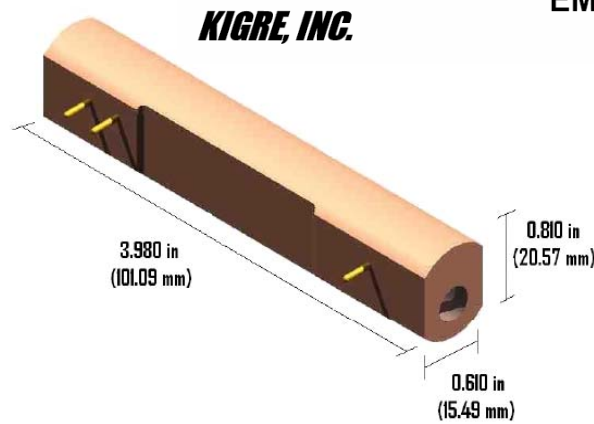
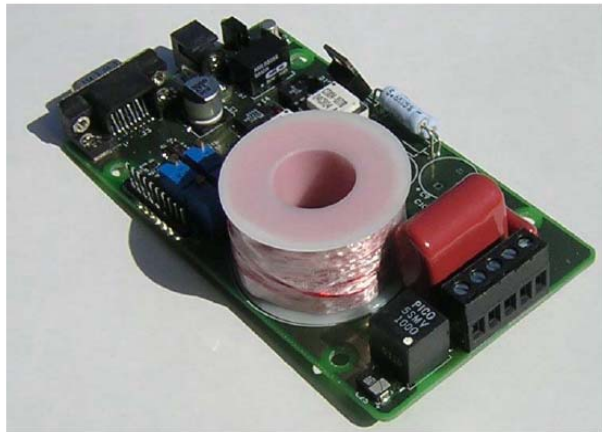
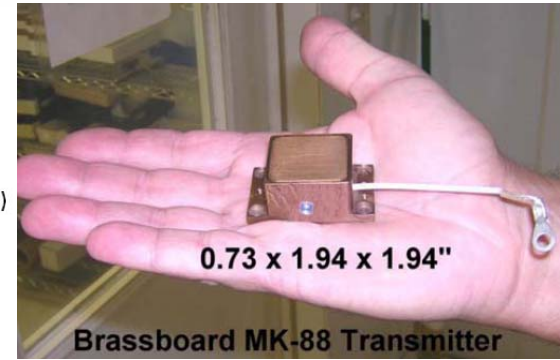


Fig. 3. LIBS spectra of two aluminum alloys.

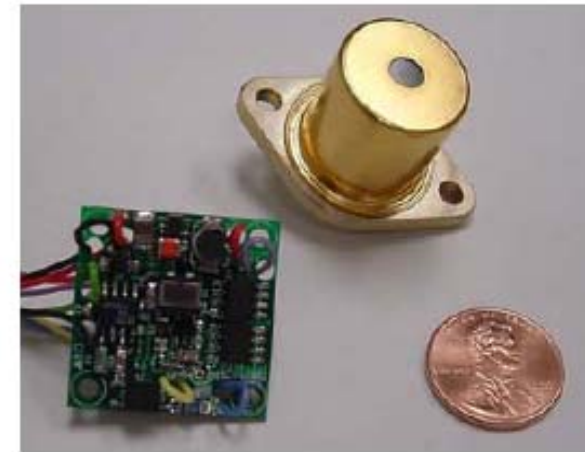
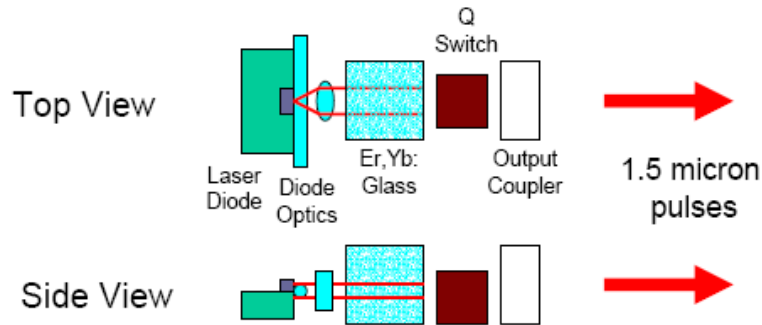


EMK-50 Er:GLASS EYE-SAFE LASER TRANSMITTER



Brassboard MK-88 Transmitter

Diode Pumped Er:Yb:glass Micro-laser MegaWatt lasers



Micro-laser in TO3 package and Diode drive circuit for use with L123 Battery.

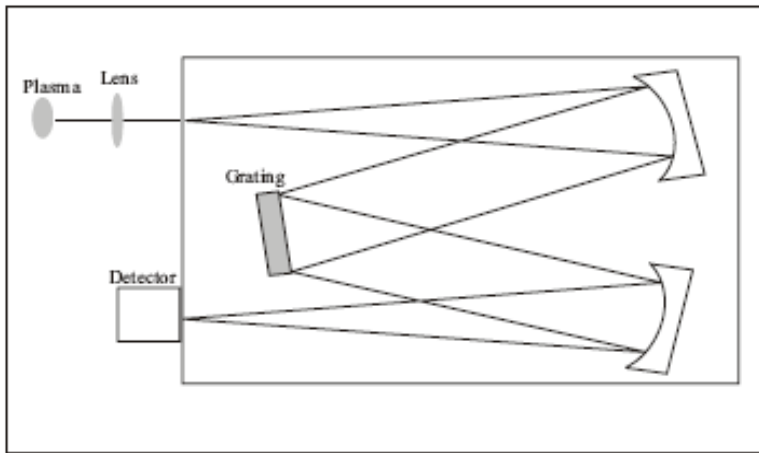
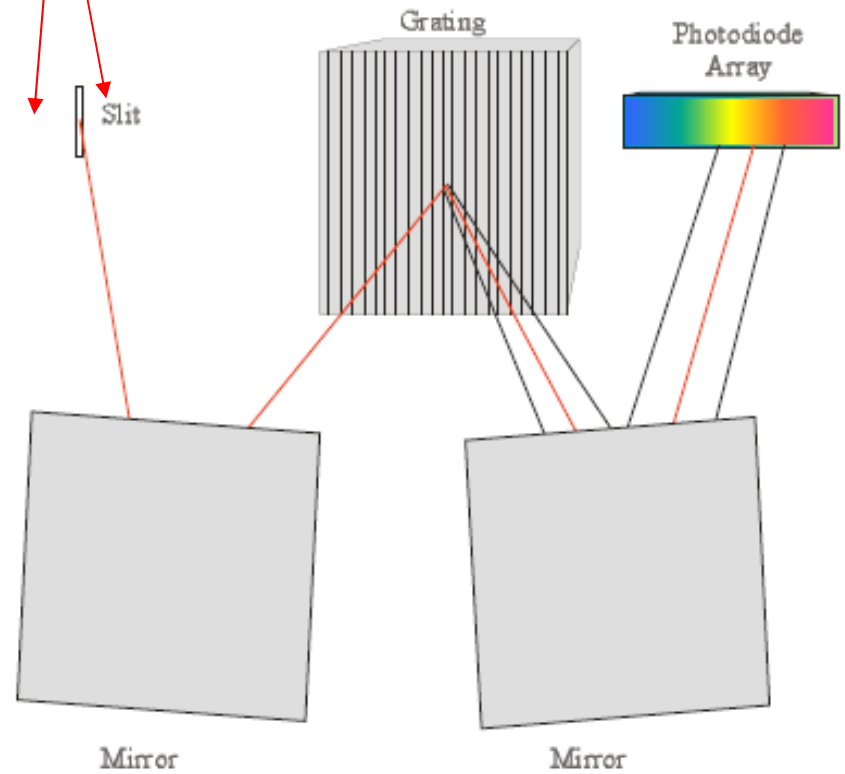


**MicroChip
NanoPulse, NanoGreen,
and NanoEyeSafe
OEM Short Package
Solid-State Lasers**

\$1K-\$10K

JDS Uniphase

Spectrometer: Czerny-Turner



Ocean-Optics USB4000-VIS-NIR Miniature Fiber Optic Spectrometer



\$3000

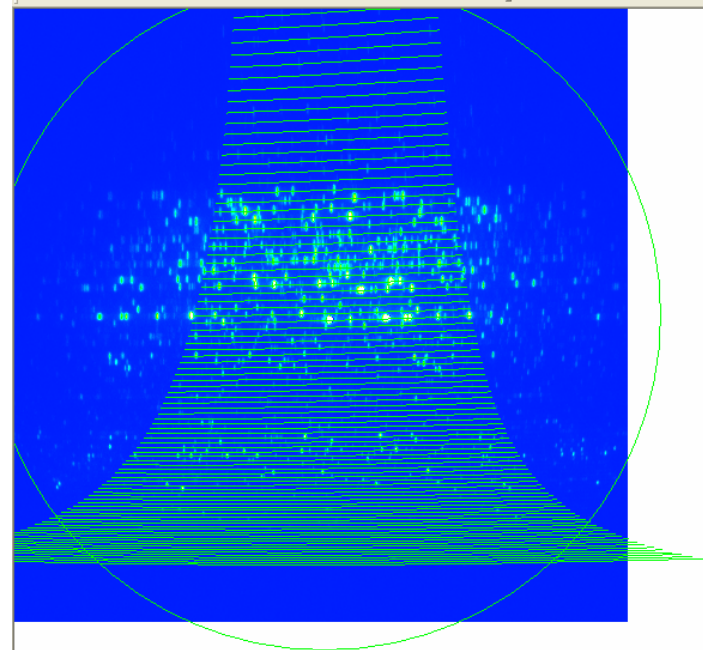
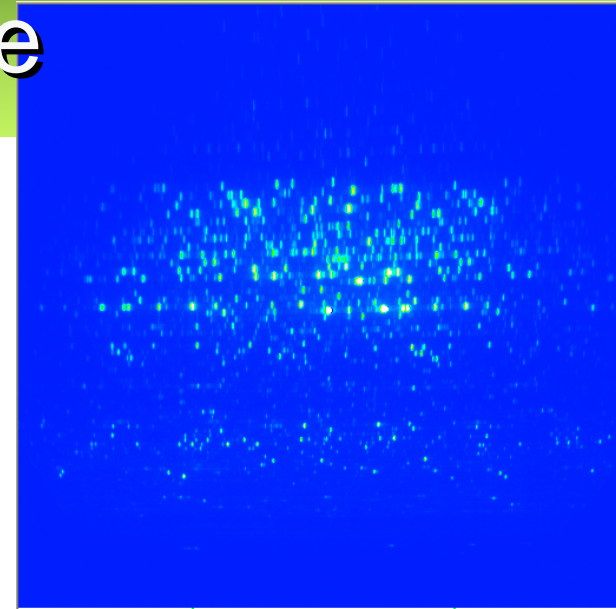
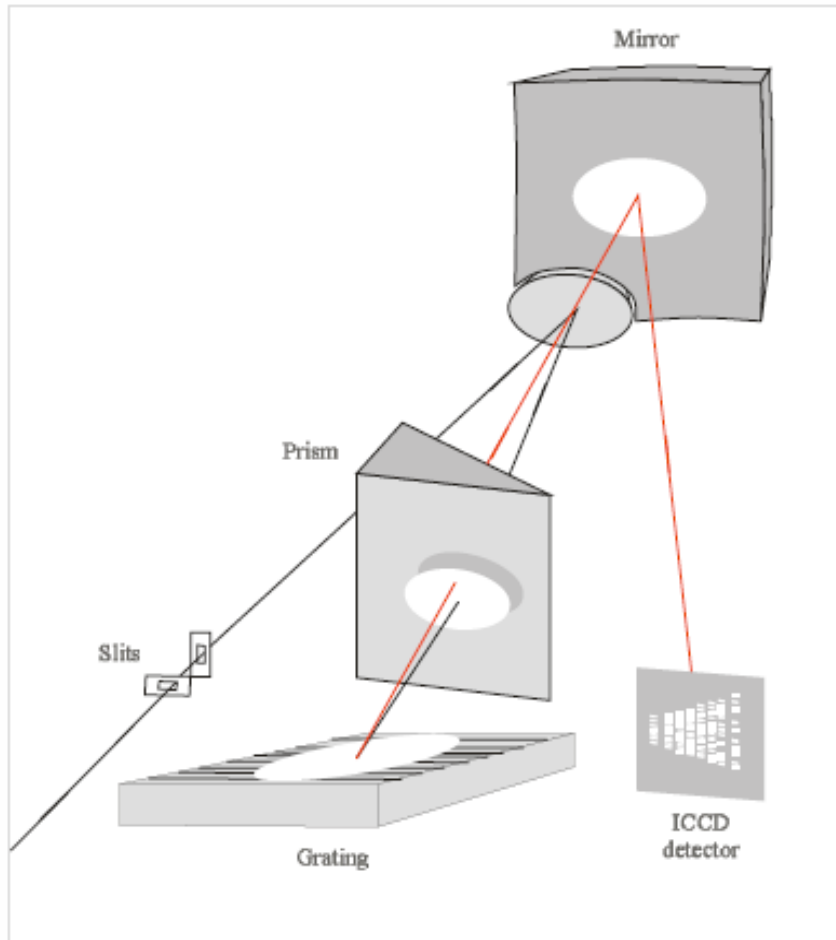


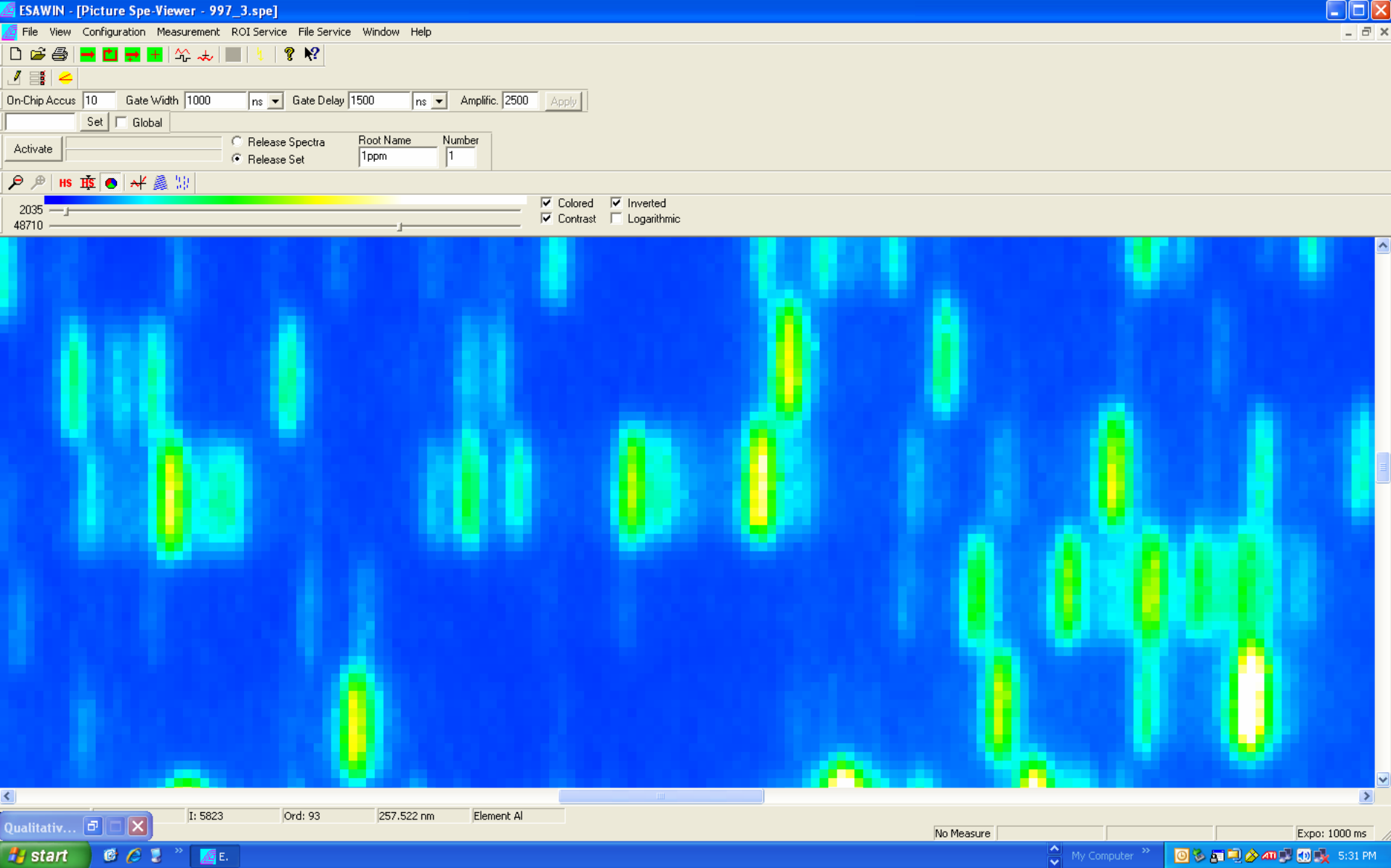
Specifications

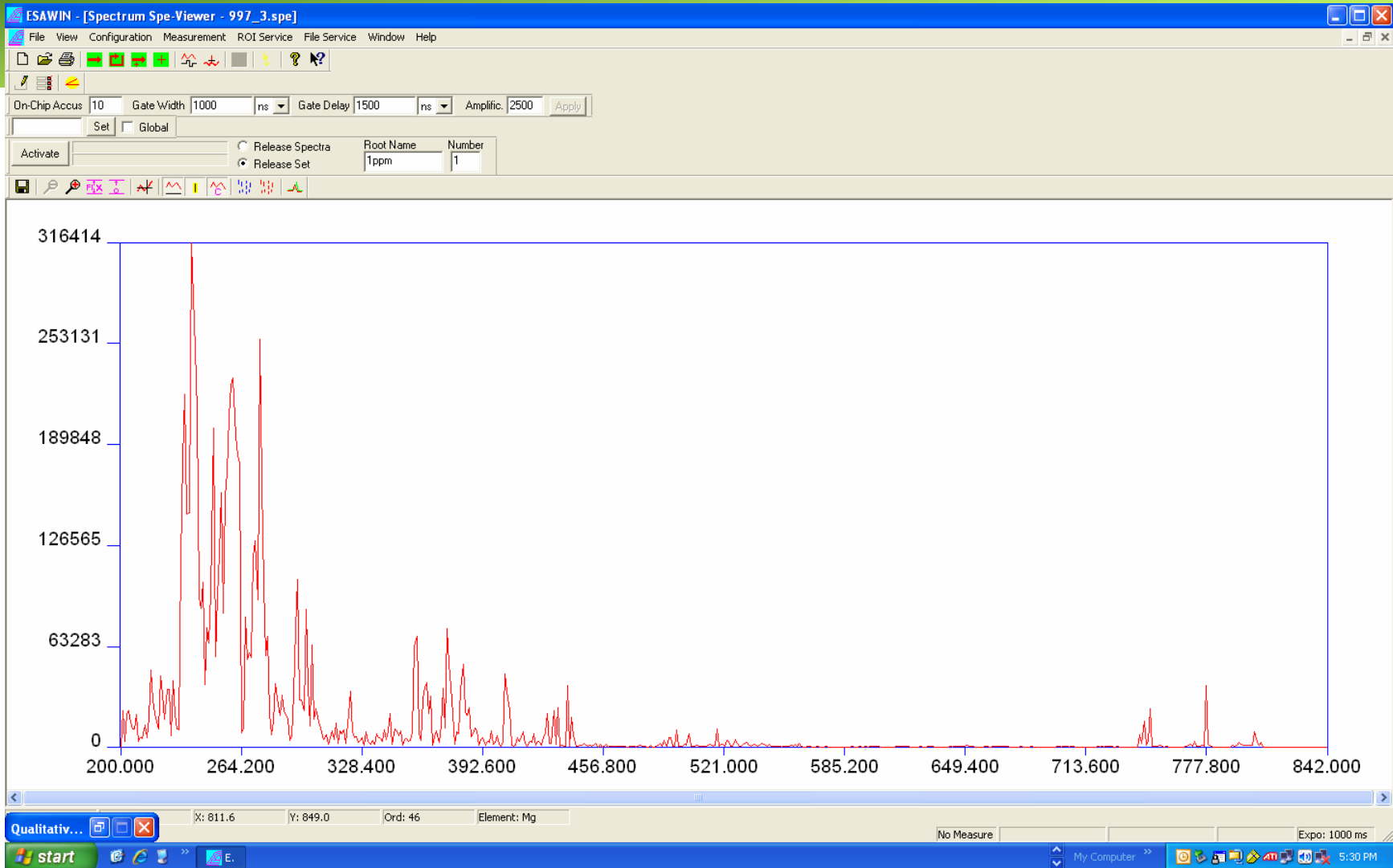
| Physical | |
|-------------------------|--|
| Dimensions: | 89.1 mm x 63.3 mm x 34.4 mm |
| Weight: | 190 grams |
| Detector Specifications | |
| Detector: | Toshiba TCD1304AP Linear CCD array |
| Detector range: | 200-1100 nm |
| Pixels: | 3648 pixels |
| Pixel size: | 8 μm x 200 μm |
| Pixel well depth: | 100,000 electrons |
| Signal-to-noise ratio: | 300:1 (at full signal) |
| A/D resolution: | 16 bit |
| Dark noise: | 50 RMS counts |
| Corrected linearity: | >99.8% |
| Sensitivity: | 130 photons/count at 400 nm; 60 photons/count at 600 nm |
| Optical Bench | |
| Design: | f/4, Asymmetrical crossed Czerny-Turner |
| Focal length: | 42 mm input; 68 mm output |
| Entrance aperture: | 25 μm wide slit |
| Grating: | Grating #3 (blazed at 500 nm) |
| OFLV filter: | OFLV-350-1000 |
| Fiber optic connector: | SMA 905 to 0.22 numerical aperture single-strand optical fiber |
| Spectroscopic | |
| Wavelength range: | 350-1000 nm |
| Optical resolution: | ~1.5 nm FWHM |
| Signal-to-noise ratio: | 300:1 (at full signal) |
| A/D resolution: | 16 bit |
| Dark noise: | 50 RMS counts |
| Integration time: | 4 ms - 10 seconds |
| Dynamic range: | 2 x 10 ⁵ (system), 1300:1 for a single acquisition |
| Stray light: | <0.05% at 600 nm; 0.10% at 435 nm |

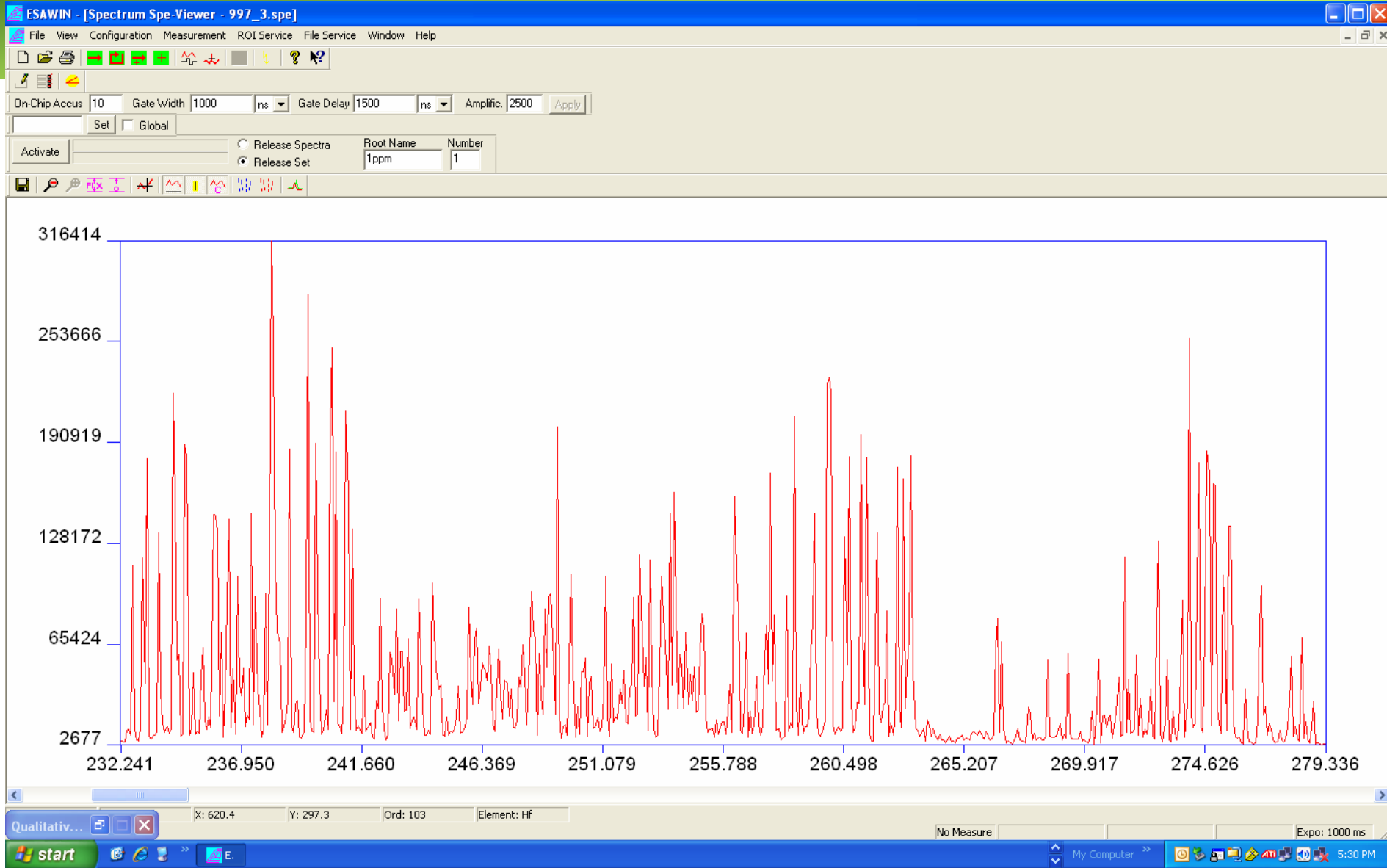
Spectrometer: Echelle

- Échelle Spectrometer









B. Salle et al., Evaluation of a compact spectrograph for in-situ and stand-off Laser-Induced Breakdown Spectroscopy analyses of geological samples on Mars missions, Spectrochimica Acta B 60, 805 – 815 (2005)

Echelle →

mini-Czerny Turner →

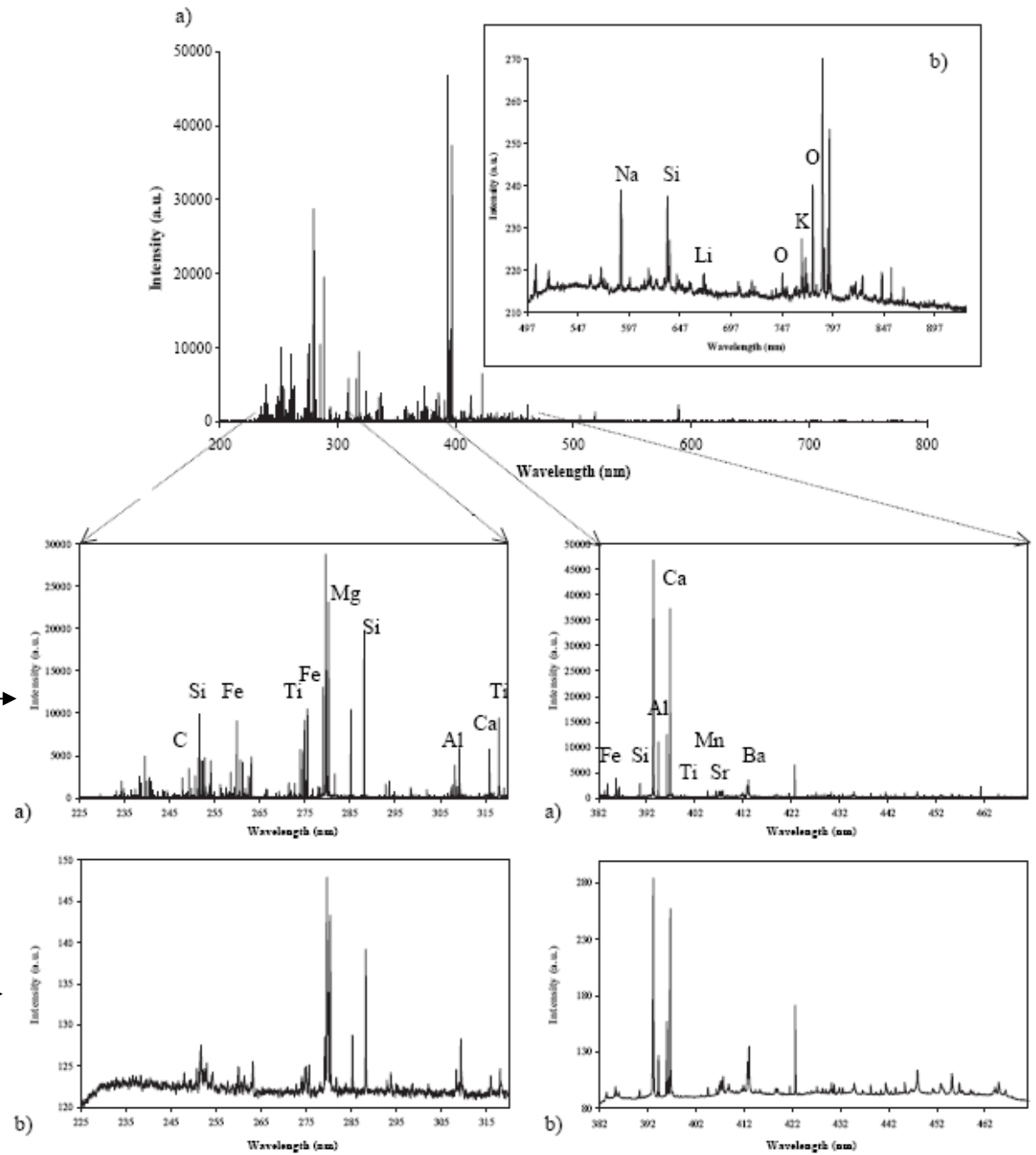


Fig. 3. Comparison of a) LLA ESA 3000 and b) Ocean Optics HR2000 spectra obtained with in-situ LIBS setup and a certified stream sediment sample.

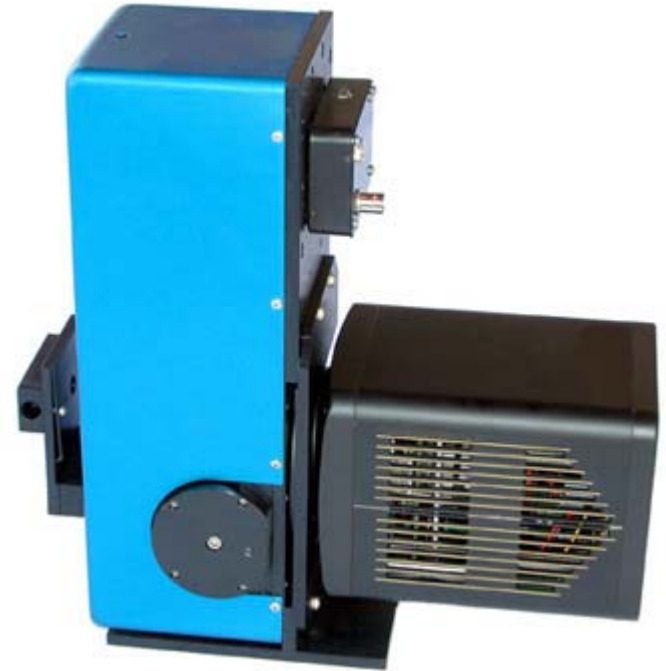


Echelle Spectrometers

LLA Instruments GmbH ESA 4000



Catalina SE200 with CCD



Andor Mechelle 5000



Aryelle 400 LaserTechnik Berlin LTB



\$35K-\$100K (Strong Euro)

Échelle References

M. Sabsabi et al., An evaluation of a commercial Échelle spectrometer with intensified charge-coupled device detector for materials analysis by laser-induced plasma spectroscopy, Spectrochim. Acta B **56**, 1011-1025 (2001)

M. Sabsabi et al., Comparative study of two new commercial Echelle spectrometers equipped with intensified CCD for analysis of laser-induced breakdown spectroscopy, Appl. Opt. **42**, 6094-6098 (2003)



ICCD

- 1-Photocathode: converts photons to electrons
- 2-Microchannel Plate (MCP): amplifies the photoelectrons
- 3-Phosphor: converts electron to photons

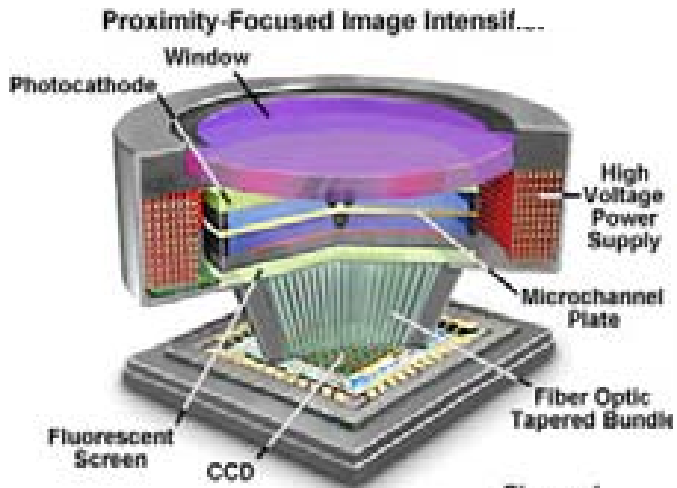
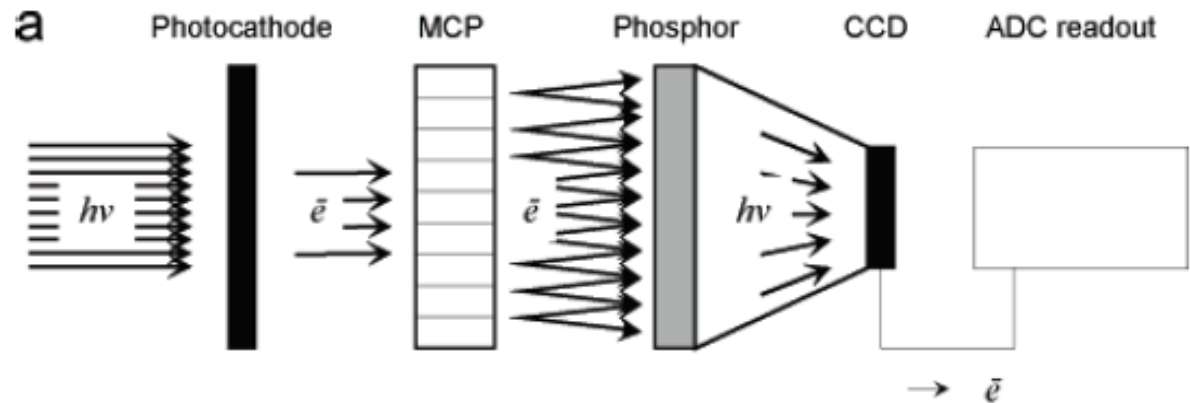
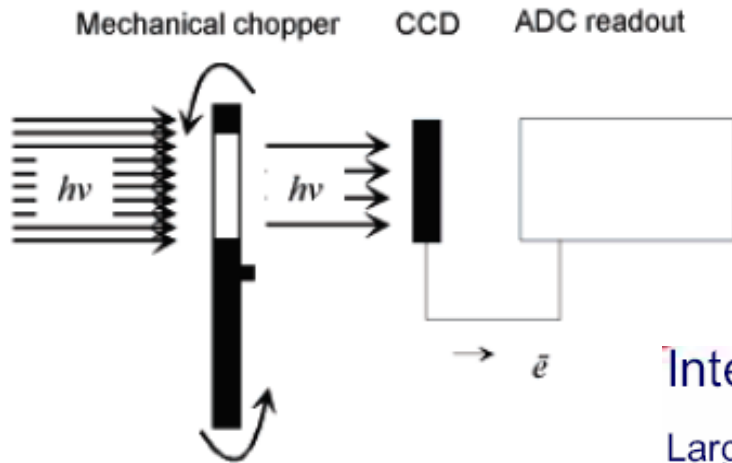
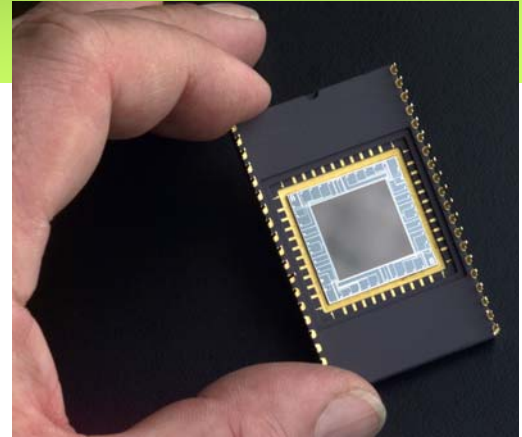


Figure 1

Poor quantum efficiency (below 20%), limited wavelength coverage, high gain, 16 bits dynamic range, simultaneous measurements, **high cost**.

Cost: 25-50 k\$

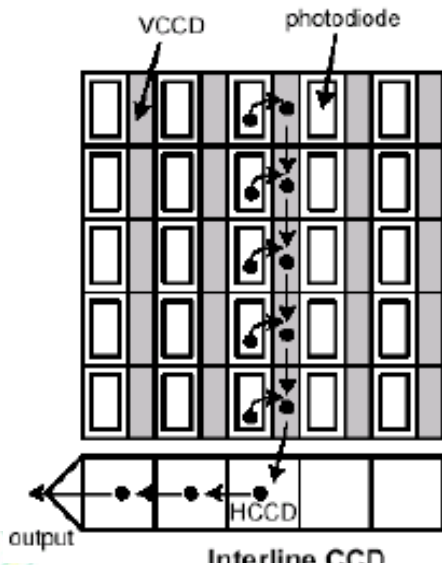
CCD



Interline CCD, Gated CCD

Larger wavelength coverage, high quantum efficiency, 12-14 bits dynamic range, **less expensive.**

Cost: 1-20 k\$



| Characteristics | Typical Range | Limitations |
|--------------------|---|---|
| Full-well capacity | 10000-500000 electrons | Defines dynamic range |
| Pixel dimensions | 6-30 μ m | Dictate spectral or spatial resolution |
| Array format | Related to pixel size and number | Dictates the active area |
| Number of pixels | 58X512 to 2048X2048 | Dictates number of resolution elements and acquisition time |
| Quantum efficiency | 0-80% | Defines ultimate sensitivity limit |
| Blooming | Present with materials having strong lines in the range 200-900nm | Cannot observe weak lines near strong lines, ghost lines |
| Read noise | Few electrons | Excess noise limits weak light detection |

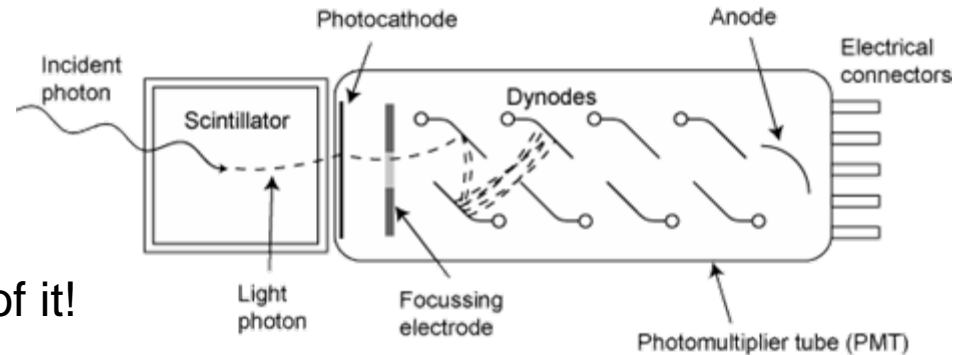
CCD References

- M. Mueller et al., Approach to Detection in Laser-Induced Breakdown Spectroscopy, Anal. Chem. **79**, 4419-4426 (2007)*
- J.E. Carranza et al., Comparison of Nonintensified and Intensified CCD Detectors for Laser-Induced Breakdown Spectroscopy, Appl. Opt. **42**, 6016-6021 (2003).*
- M. Sabsabi et al., Critical evaluation of gated CCD detectors for laser-induced breakdown spectroscopy analysis, Spectrochim. Acta B **60**, 1211-1216 (2005)*

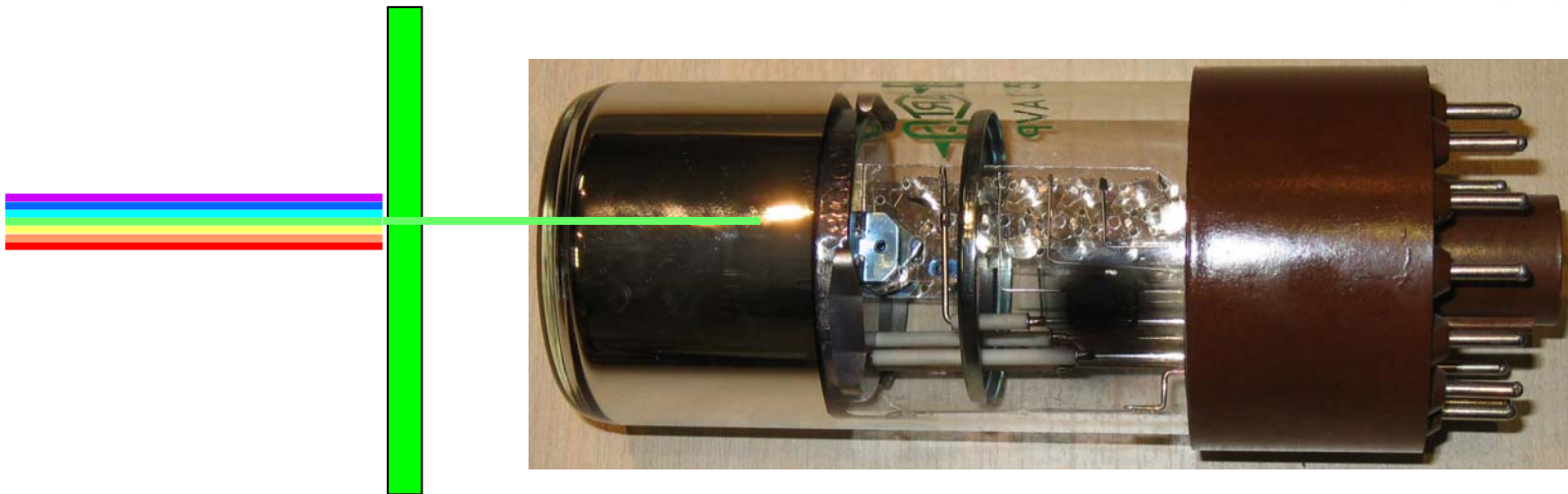


photomultiplier tube

- You may not need a spectrometer! (adds complexity, cost)



put an interference filter in front of it!



Monochromatic, low quantum efficiency, bulky, larger wavelength coverage, high dynamic range, **cheap**.

Cost: 1-3 k\$

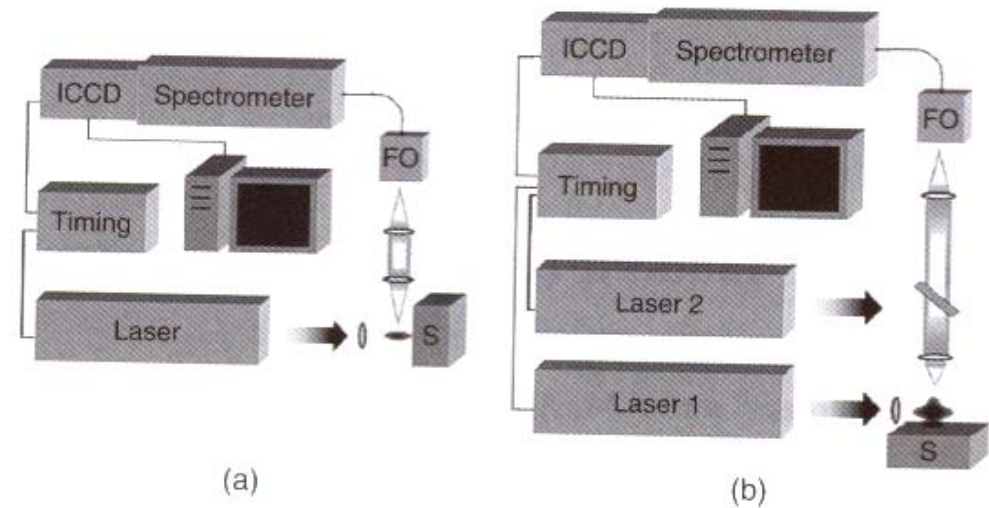
Outline

1. Introduction to LIBS
2. Physics of the plasma formation and observation
3. Instrumentation
4. **Advanced techniques**



Dual-pulse LIBS

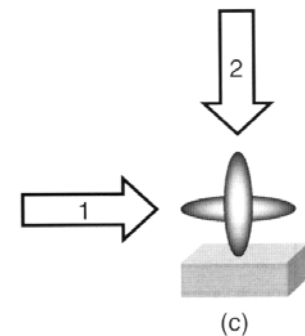
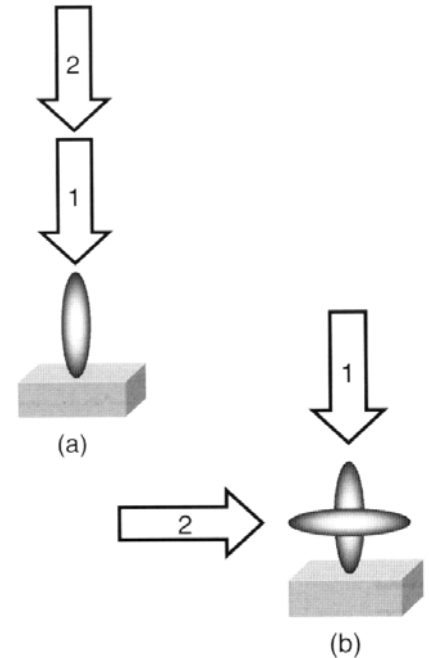
- The use of two non-coincident laser pulses increases emission intensity and SNR
- Variables to be determined
 - pulse geometry
 - pulse order
 - inter-pulse timing
 - energy of pulses



Dual-pulse pulse geometry

J. Scaffidi et al. "Dual pulse LIBS," in Laser-Induced Breakdown Spectroscopy (2007) pp. 137-150

- (a) collinear. First and second pulses are both focused onto or into the sample
- (b) orthogonal reheating. A single ablative pulse (1) is followed by a post-ablative reheating pulse (2) focused into the plasma plume.
- (c) orthogonal pre-ablative spark. A pre-ablative air spark (1) focused up to several mm above the sample surface is followed by a single ablative pulse (2) focused onto or into the sample.



Dual-pulse enhancement

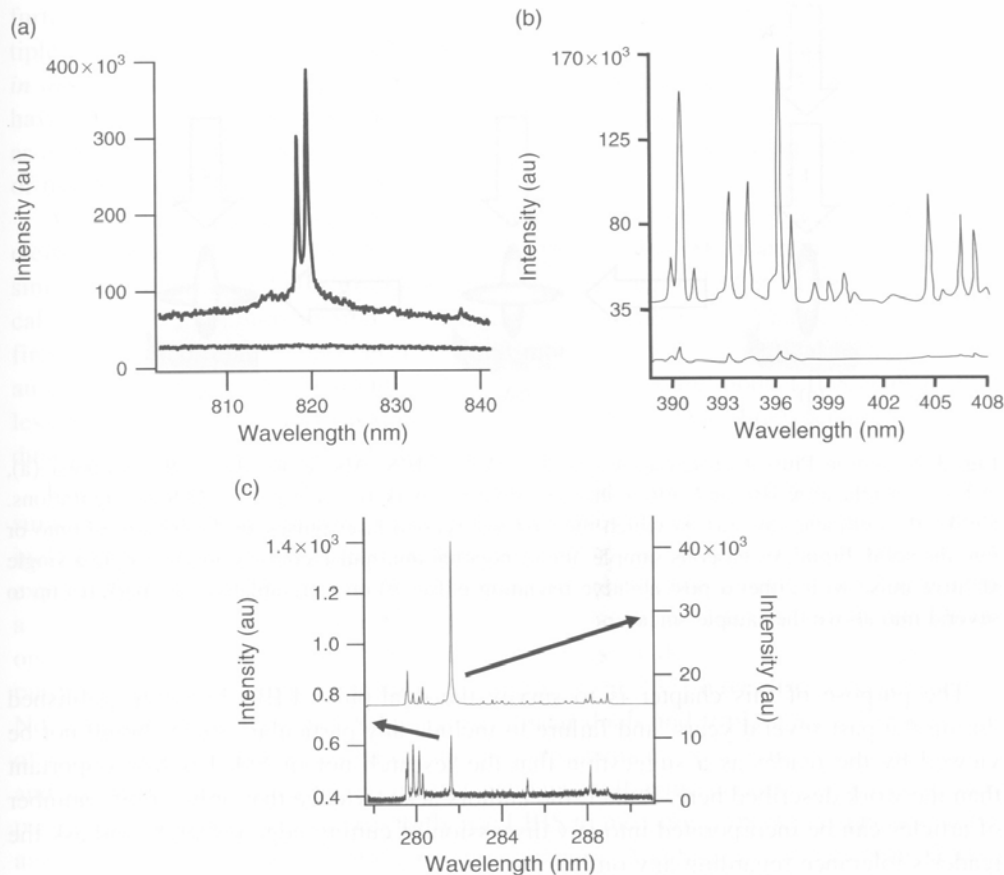
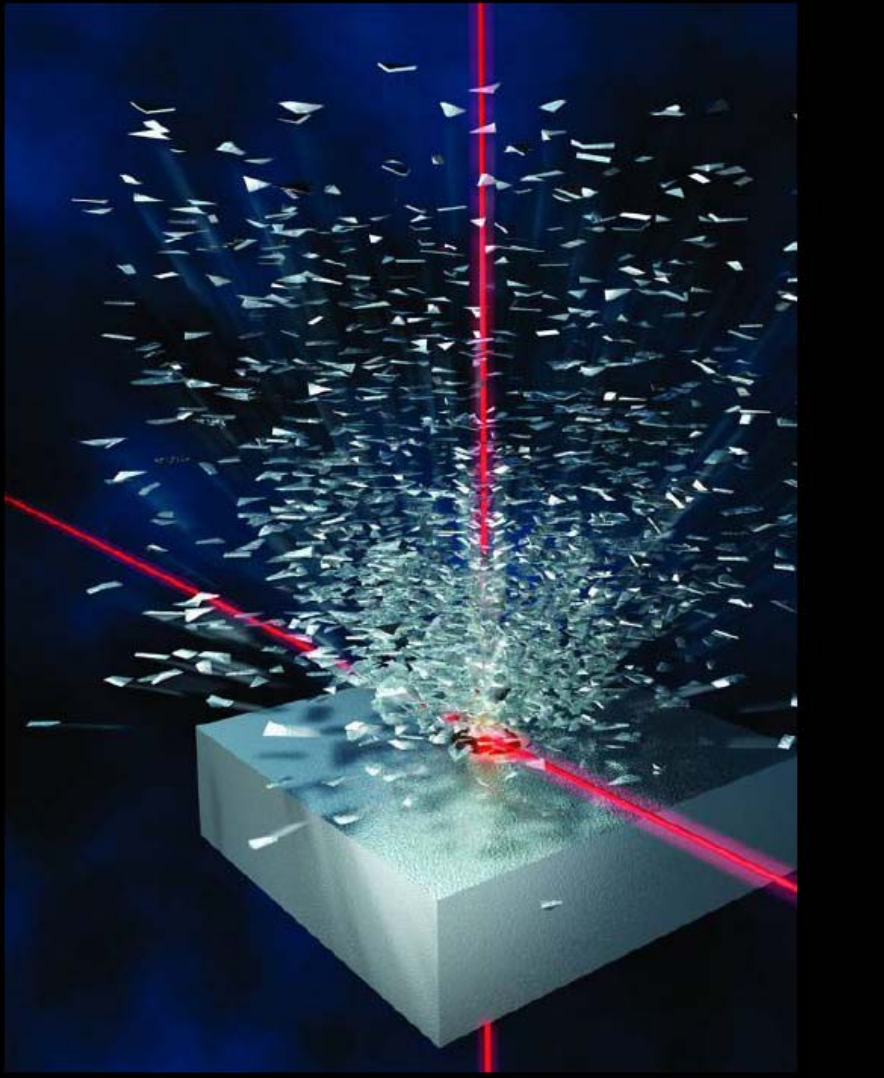


Fig. 2. Dual-Pulse LIBS Atomic Emission and Signal-to-Noise Enhancements. Above are shown sample spectra from Angel, et. al. (a, b) and Sabsabi, et. al. (c) showing the neutral atomic (a, b) and ionic (c) emission and signal-to-noise enhancements possible with orthogonal pre-ablative spark (a, b) and collinear (c) dual-pulse LIBS (top spectra) of dissolved sodium in aqueous solution (a), vitrified glass simulants in air (b), and an aluminum standard in air (c). The lower traces are the corresponding ns single-pulse LIBS (bottom spectra).

- Different pulse configurations suggest the existence of different sources for dual-pulse LIBS enhancements
- Energetic coupling between first LIP and second pulse (reheating/collinear)
- sample heating (pre-ablative)
- reductions in atmospheric pressure or number density effects

Dual-pulse enhancement



Emission Enhancement Mechanisms in Dual-Pulse LIBS

Remote analysis for challenging applications may be possible because LIBS uses only light and collects only photons.

Jon Scaffidi

University of Pittsburgh

S. Michael Angel

University of South Carolina

David A. Cremers

Analytical chemistry **78**, p. 24.

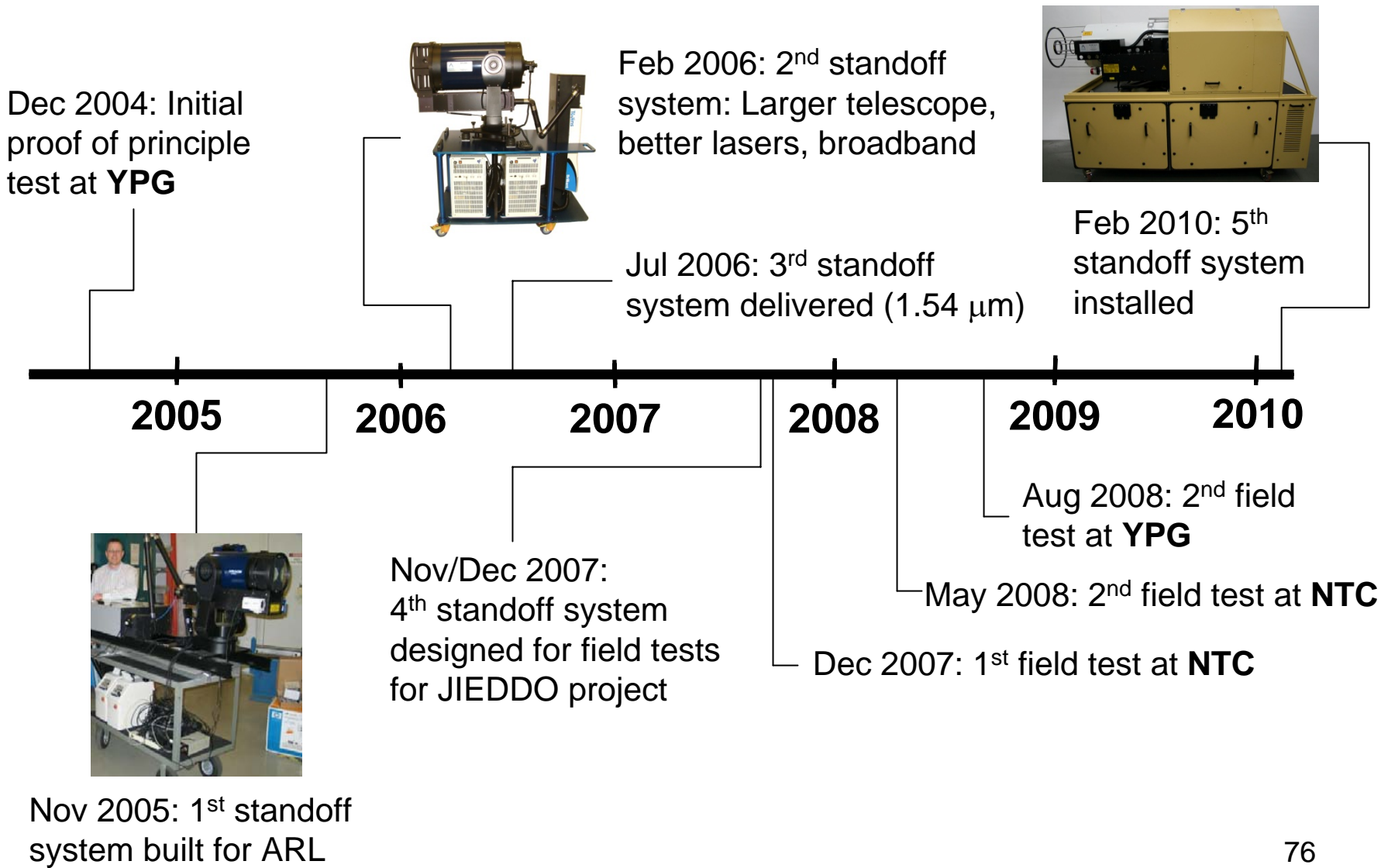
Applied Research Associates 74



Stand-off LIBS / Remote LIBS

- Not the same
 - Stand-off: the laser is delivered through space to a target. (the target is physically separated, typically 1-100 m, from the laser and light collection optics)
 - Remote: the laser and light collection optics are in physical proximity (contact) to the target but they are remote from the user/operator.







Video Camera

Telescope

Laser Head

INVITED PAPER

Filamentation “remote” sensing of chemical and biological agents/pollutants using only one femtosecond laser source

S.L. Chin · H.L. Xu · Q. Luo · F. Th  berge · W. Liu · J.F. Daigle · Y. Kamali ·
P.T. Simard · J. Bernhardt · S.A. Hosseini · M. Sharifi · G. M  jean · A. Azarm ·
C. Marceau · O. Kosareva · V.P. Kandidov · N. Ak  zbek · A. Becker · G. Roy ·
P. Mathieu · J.R. Simard · M. Ch  teauneuf · J. Dubois

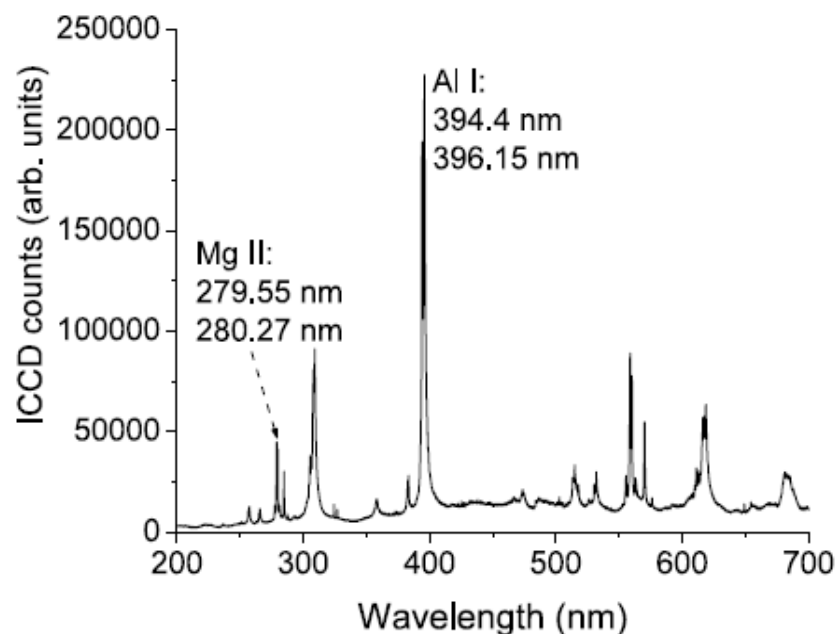


Fig. 9 Filament-induced breakdown spectrum taken for aluminum sample located 50 m away. The laser pulse energy was 108 mJ, and the ICCD gate width was 2 μ s with a time delay of $t = -33$ ns with respect to the laser pulse arriving time on the target ($t = 0$)



a

Filament-induced remote surface ablation for long range laser-induced breakdown spectroscopy operation ☆

Ph. Rohwetter^a, K. Stelmaszczyk^a, L. Wöste^a, R. Ackermann^b, G. Méjean^b, E. Salmon^b,
J. Kasparian^b, J. Yu^{b,*}, J.-P. Wolf^b

Spectrochimica Acta Part B 60 (2005) 1025 – 1033

(b)

(a)

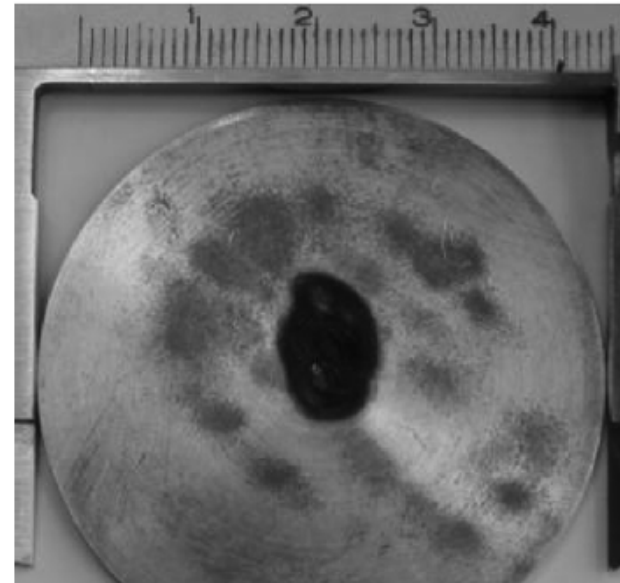
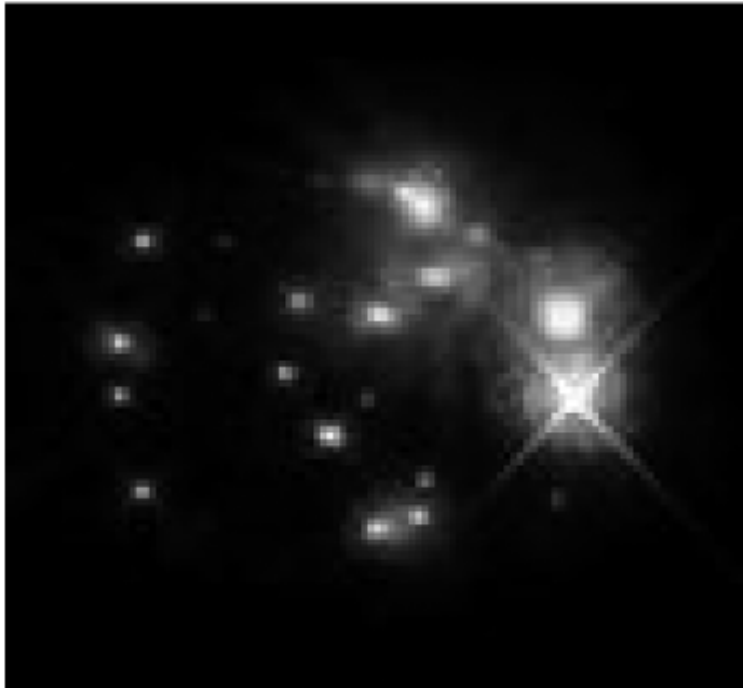


Fig. 1. (a) Picture of a typical laser beam profile with many filaments (single shot picture). On the picture, filaments appear as bright spots. Notice that each filament is surrounded by a conical emission in the visible range that increases the apparent size of the filaments on the picture to the order of 10^2 . The whole beam had a diameter of 5 cm. (b) Picture of a copper sample irradiated by a large number of laser pulses. The black spot in the middle is due to irradiation with a focused beam from a distance of 25 m. The grey spots are results of filament irradiation of the sample placed at a distance of 90 m. Random walks of a filament around its mean position due to inhomogeneities in air lead to a superficial ablation on a surface much larger than filament diameter. The scale of this picture is indicated by a ruler with a 1 cm graduation in the upper part of the picture.

Filament-induced remote surface ablation for long range laser-induced breakdown spectroscopy operation[☆]

Ph. Rohwetter^a, K. Stelmaszczyk^a, L. Wöste^a, R. Ackermann^b, G. Méjean^b, E. Salmon^b,
J. Kasparian^b, J. Yu^{b,*}, J.-P. Wolf^b

Spectrochimica Acta Part B 60 (2005) 1025 – 1033

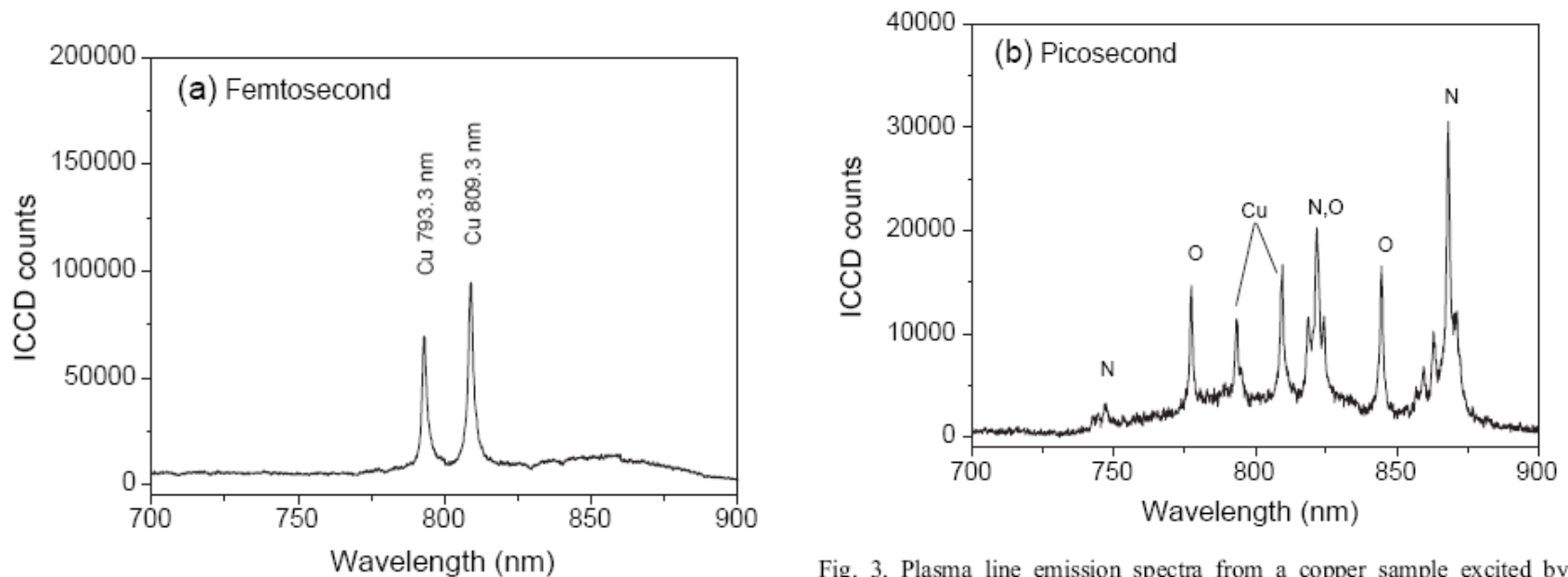
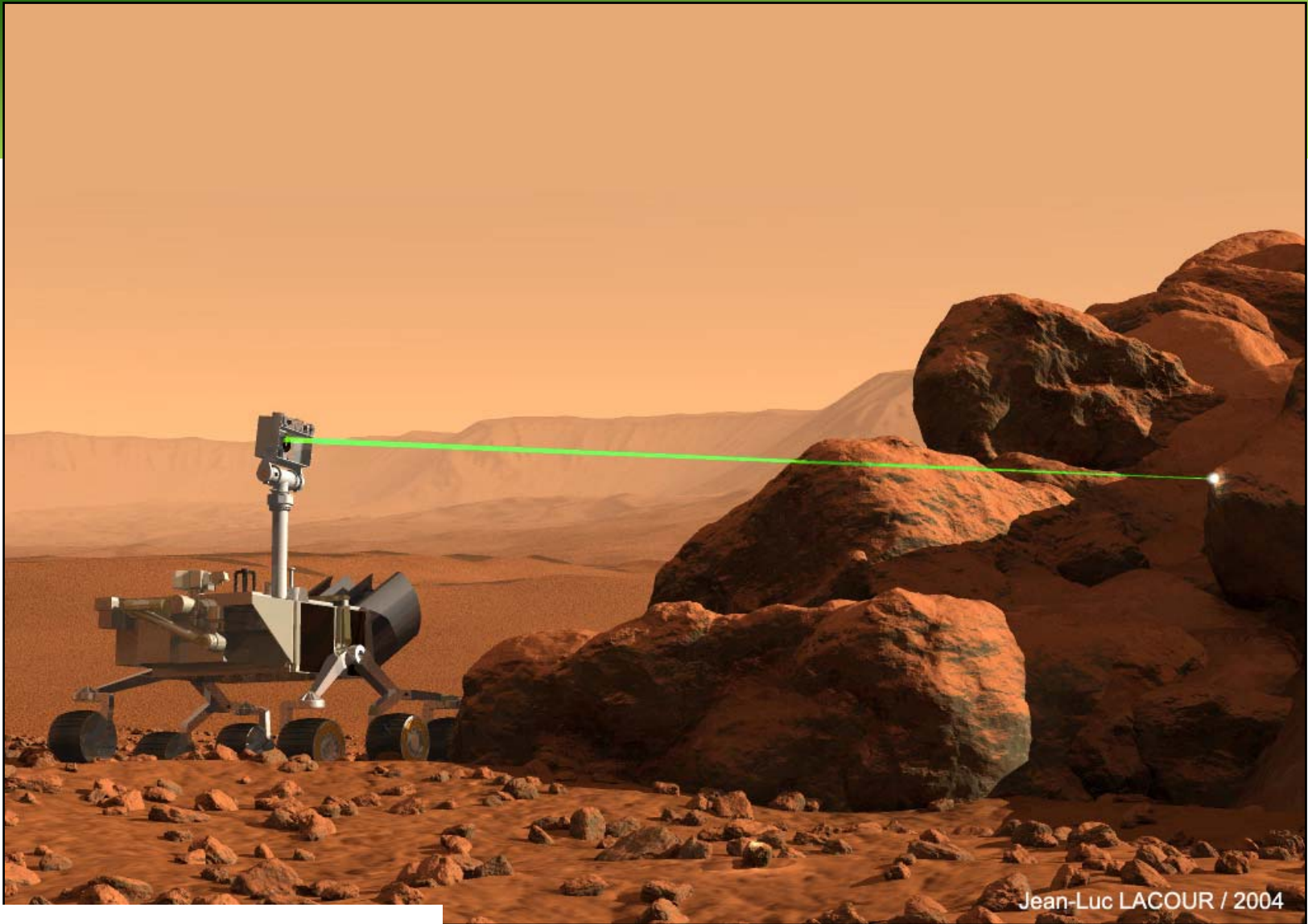


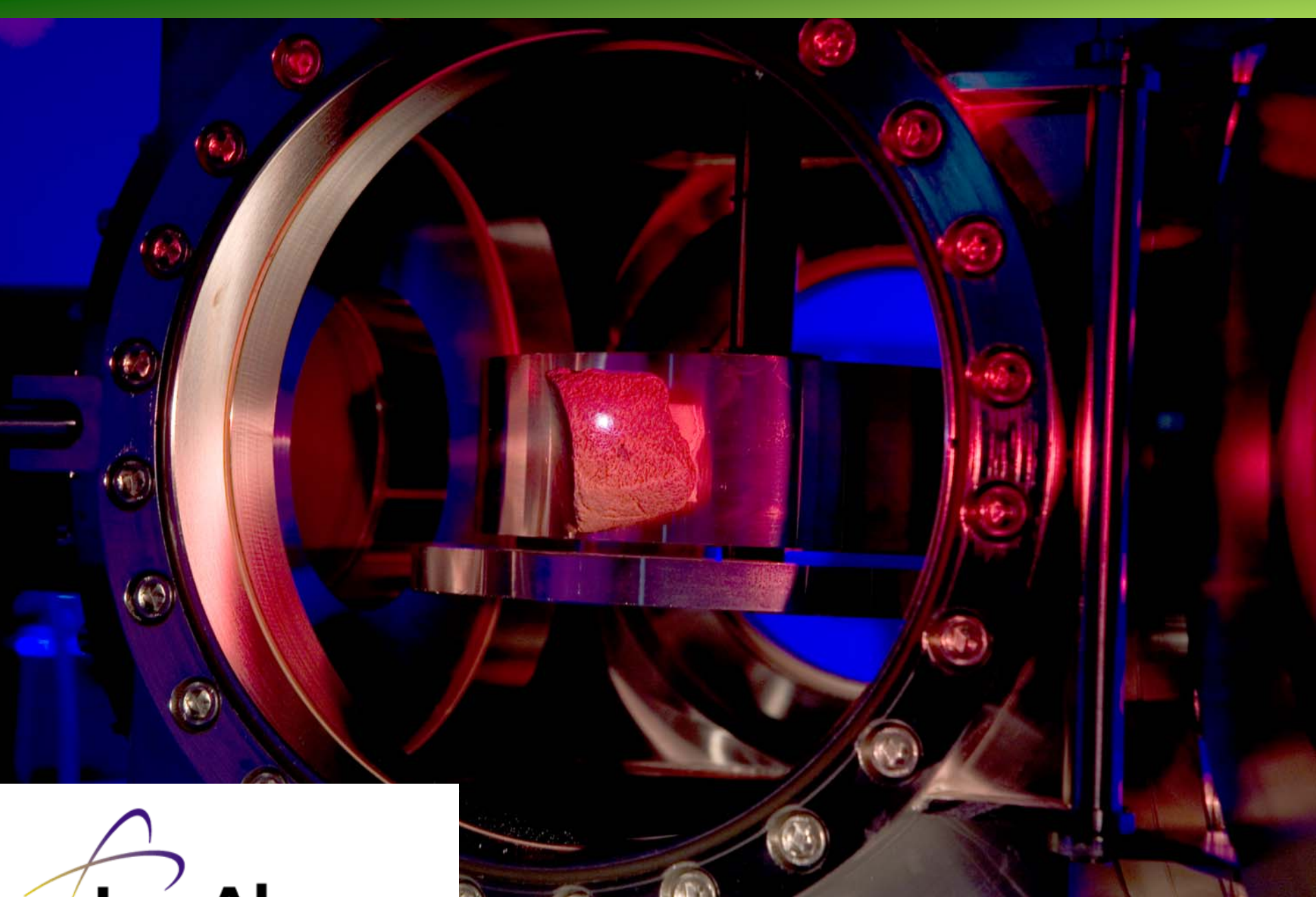
Fig. 3. Plasma line emission spectra from a copper sample excited by femtosecond pulses (a), and by picosecond pulses (b). With femtosecond pulses, two copper atomic lines in the spectral range from 700 nm to 900 nm are observed without any background lines. In contrast, with picosecond (or nanosecond pulses), the copper lines are mixed with background lines from ambient air: atomic oxygen (O) and atomic nitrogen (N) lines.





DANGER
INVISIBLE LASER RADIATION -
AVOID DIRECT EXPOSURE TO BEAM
Nd:KGW; Diode
1067 nm; 785 nm
4 ns / 24 mJ; 50 mW
CLASS IIIb Laser Product

CHEMCAM MAST UNIT
EM MODEL
REMOVE BEFORE TEST



Chemometrics

- Chemometrics is advanced signal-processing that statistically (mathematically) identifies similarities and differences in LIBS spectra.
- Can express the spectra in a basis set that maximizes the differences for real-time classification and/or identification



Chemometrics

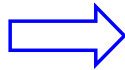
- Intensity of lines, ratios of intensities used in a statistical multi-variate analysis
 - discriminant function analysis (DFA)
 - principal component analysis (PCA)
 - partial least squares – discriminant analysis (PLS-DA)
 - linear discriminant analysis (LDA)



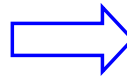
Chemometrics

How a DFA is Performed

(1) The intensities of 13 emission lines are normalized by the sum of all intensities and separated into N groups



(2) “**Canonical Discriminant Functions**” are constructed from arrays of groups. For discrimination between N groups, $N-1$ canonical discriminant functions are constructed

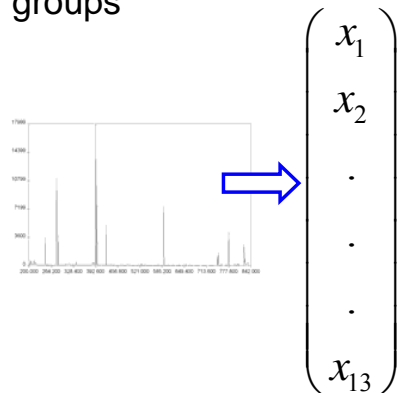


(3) For each spectrum, $N-1$ **discriminant function scores** are calculated

$$DF^j = b_0^j + \sum_{k=1}^{13} b_k^j x_k$$



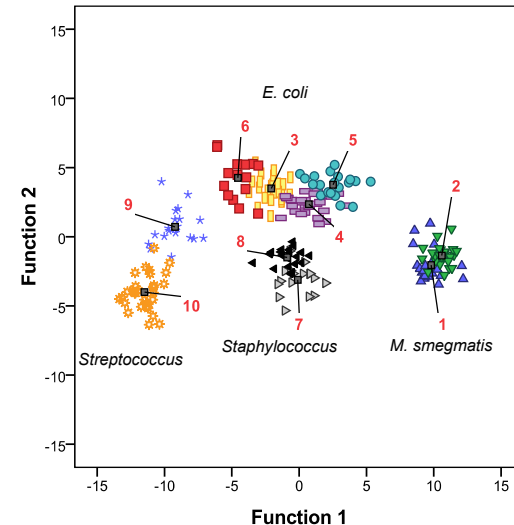
(4) Each spectrum is plotted according to its calculated discriminant function scores as a data point on a graph (often 2D).



$$(b_0^1 \quad b_1^1 \quad b_2^1 \quad \cdot \quad \cdot \quad b_{13}^1)$$

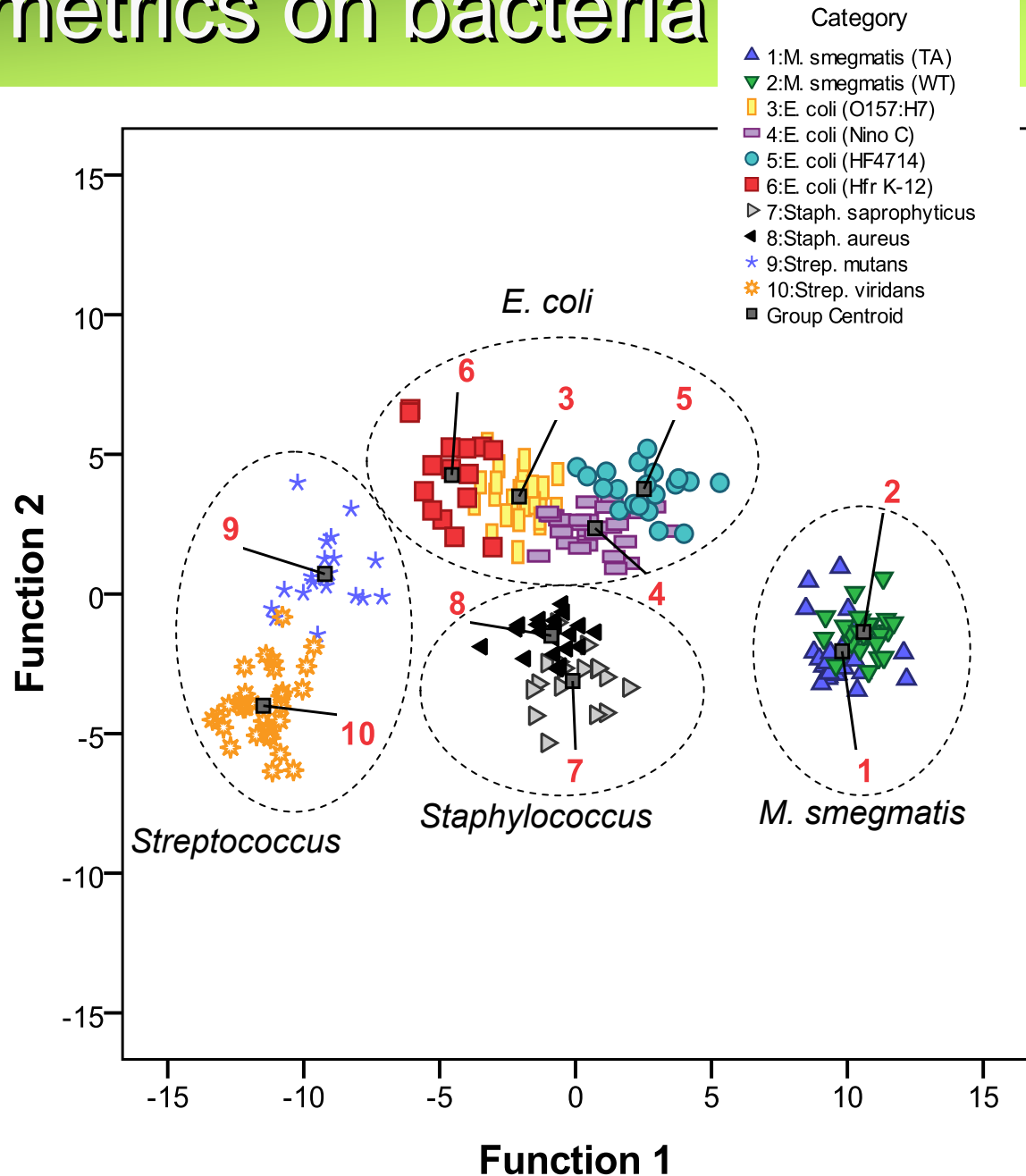


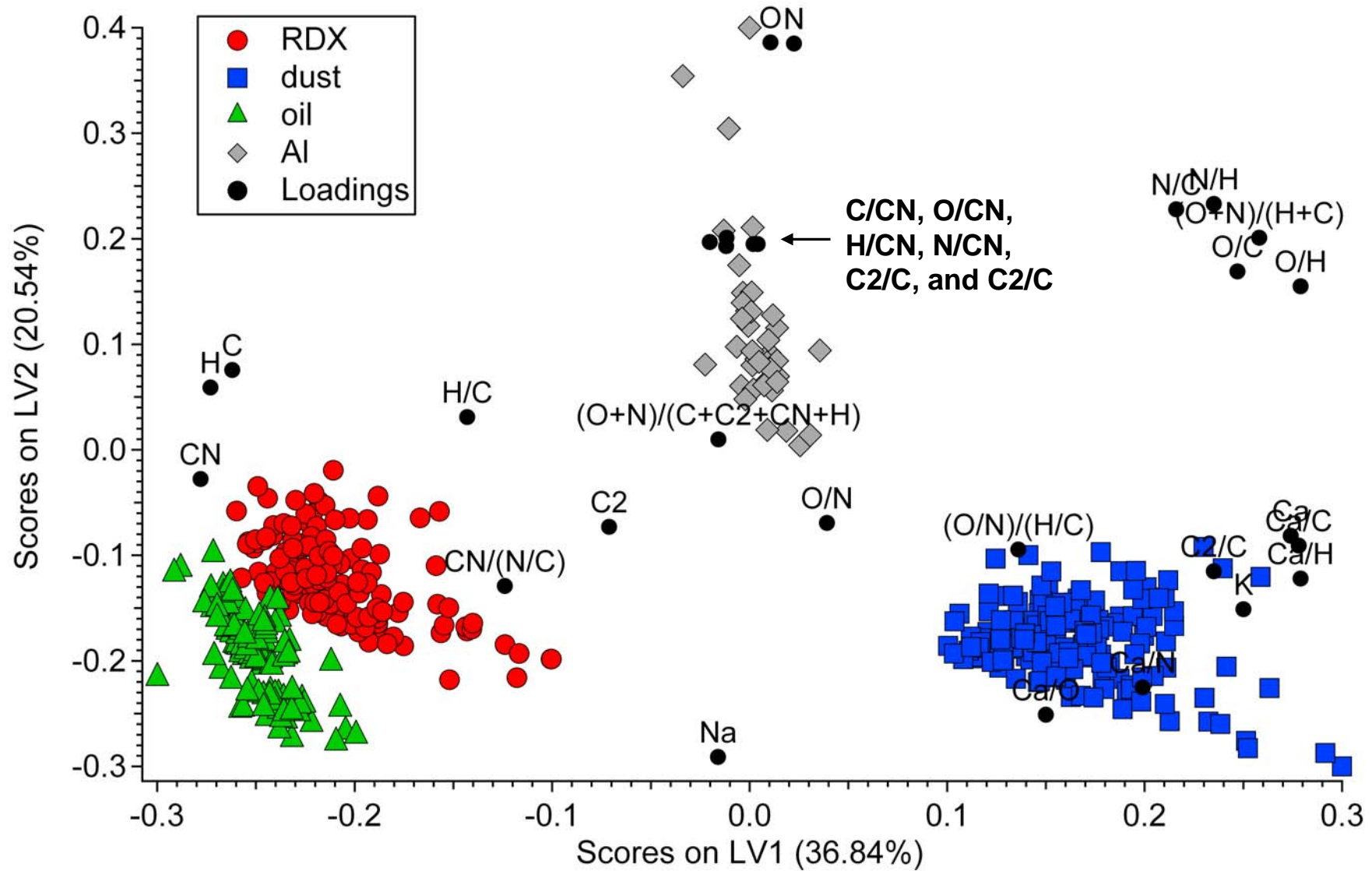
$$(b_0^{N-1} \quad b_1^{N-1} \quad b_2^{N-1} \quad \cdot \quad \cdot \quad b_{13}^{N-1})$$

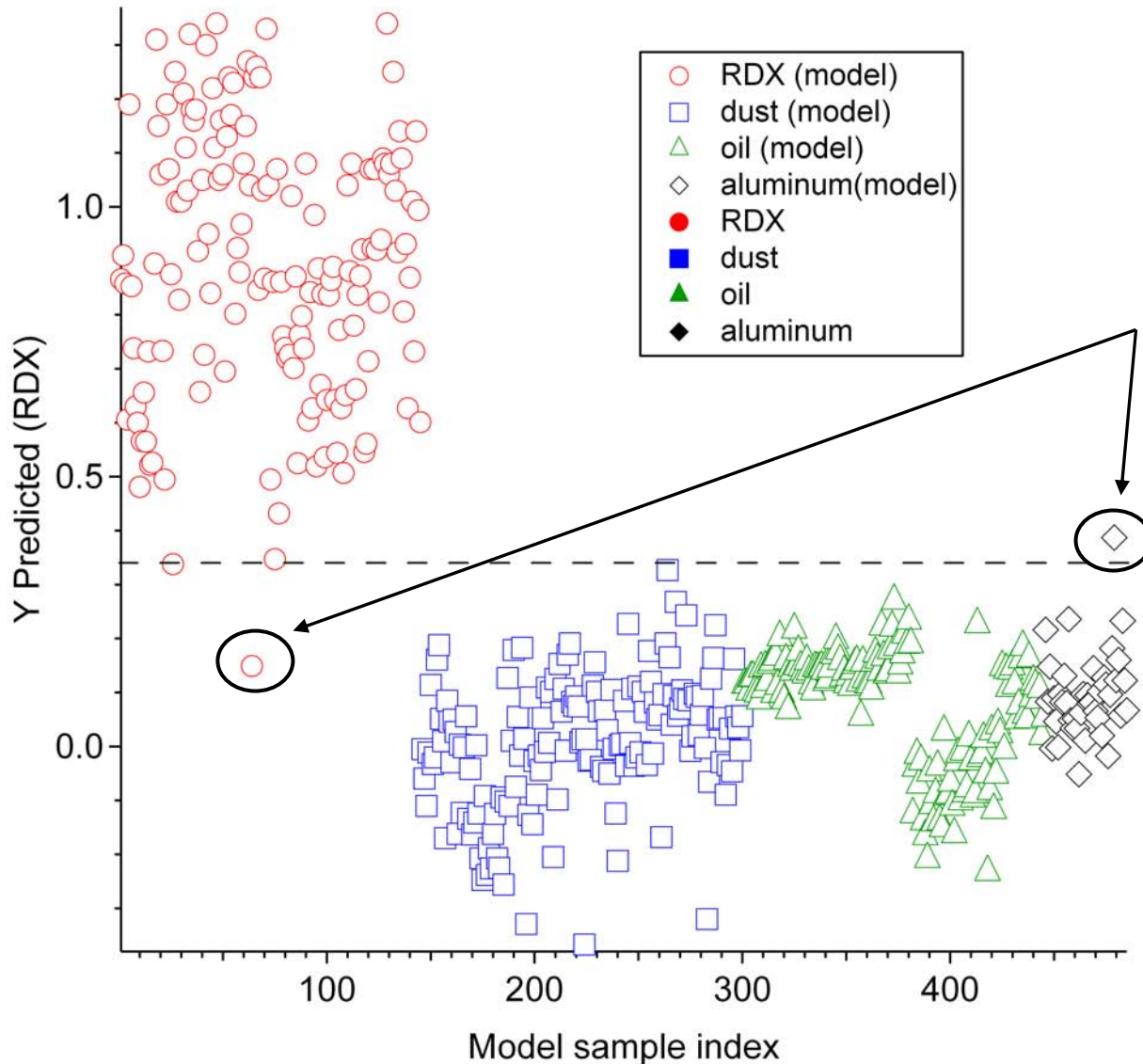


Chemometrics on bacteria

- Each data point is an entire spectrum
- Discriminant function analysis (DFA) performed on 13 LIBS emission lines
- Software was SPSS v18
- Program was only told there were 10 groups of bacteria
- Identification of an unknown sample can be done using this pre-compiled library





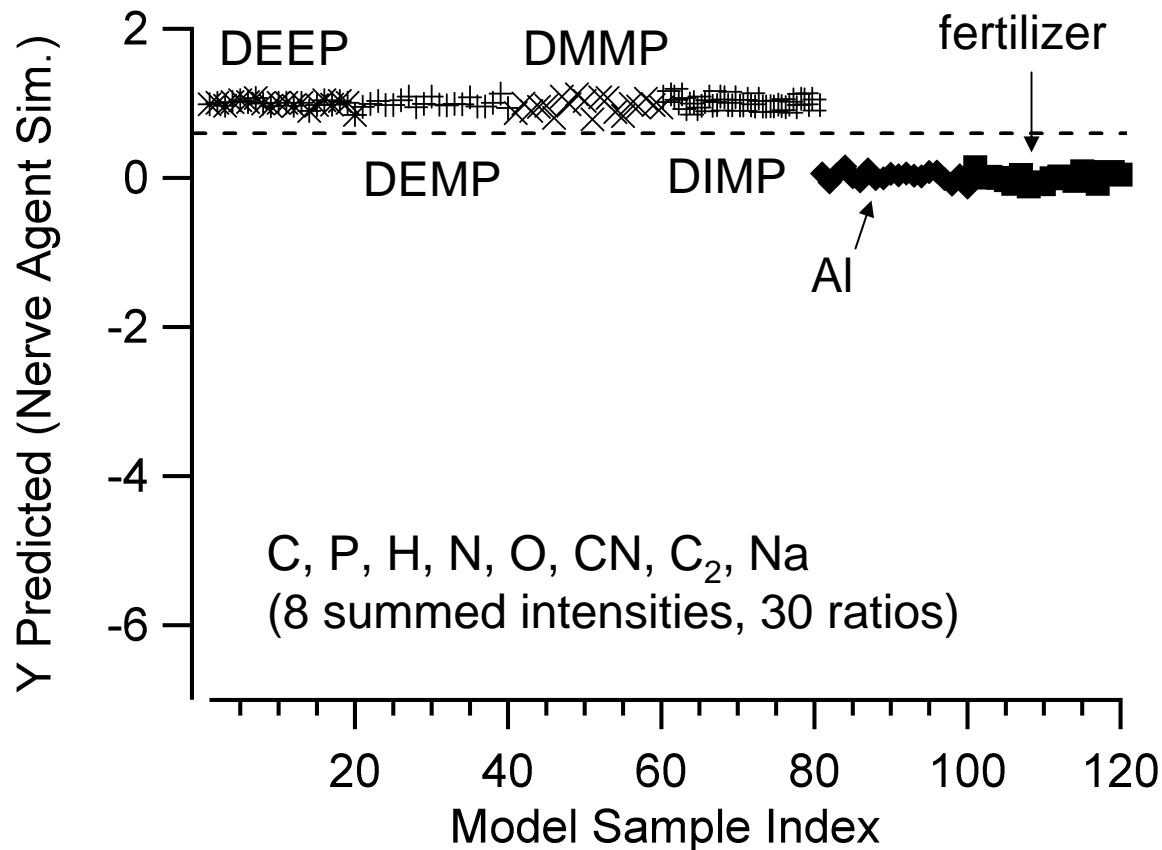


Removing statistical outliers from the model when using single-shot spectra decreases the robustness of the model!

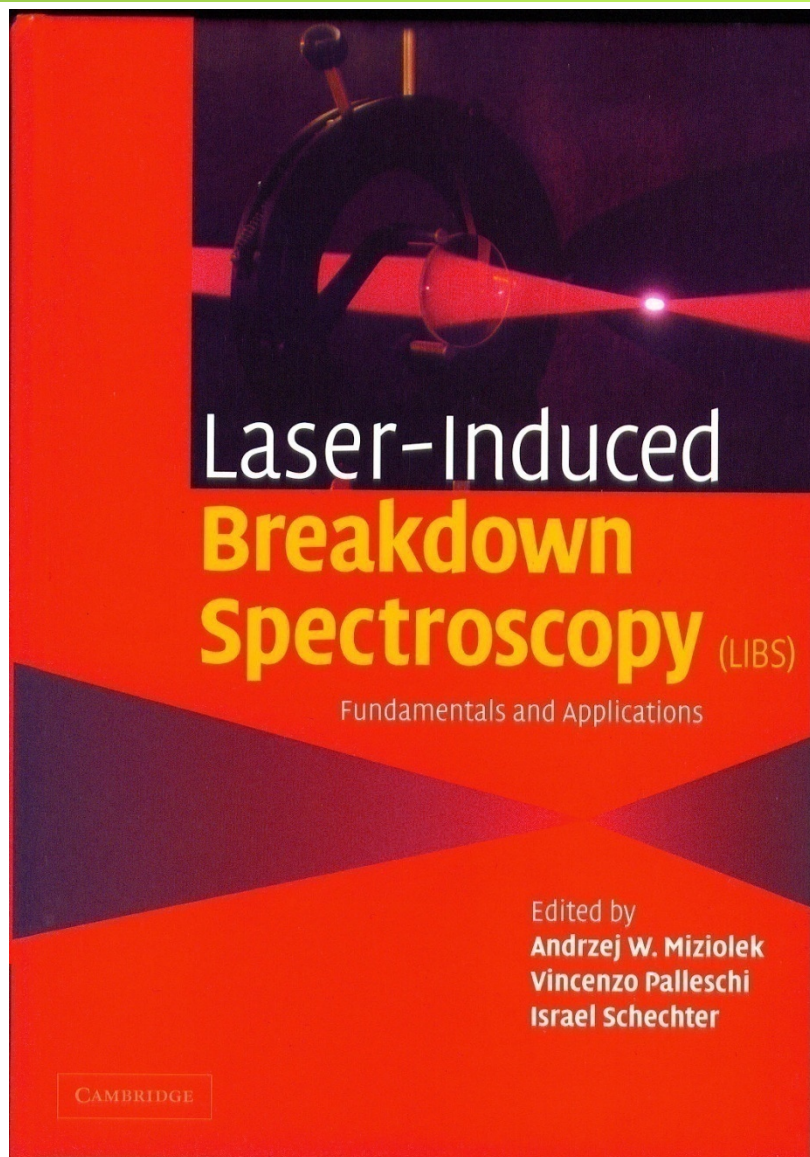
- decreased sensitivity and selectivity

100% true positives with no false negatives

(Model includes data at 20 meters and 30 meters)



- TEP was correctly identified as a nerve agent simulant (0% false negatives) while none of the interferents registered as false positives



Laser Induced Breakdown Spectroscopy (LIBS)

Edited by-

Andrzej W. Miziolek

U.S. Army Research Laboratory, USA

Vincenzo Palleschi

Istituto per I Processi Chimico-Fisici, Italy

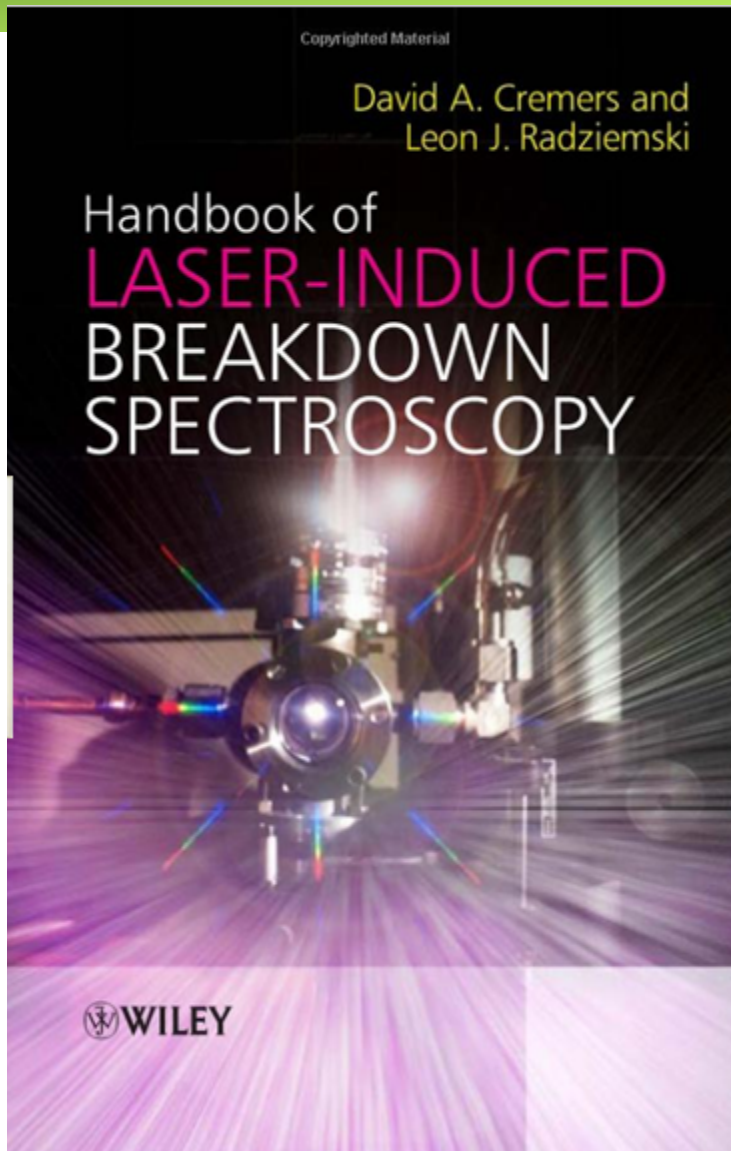
Israel Schechter

Technion - Israel Institute of Technology, Haifa, Israel

Hardback (ISBN-13: 9780521852746 | ISBN-10: 0521852749)

Published September 2006 | 638 pages | 247 x 174 mm

Cambridge University Press



Handbook of Laser-Induced Breakdown Spectroscopy

Edited by-

David A. Cremers

*Applied Research Associates, Inc., Albuquerque,
NM*

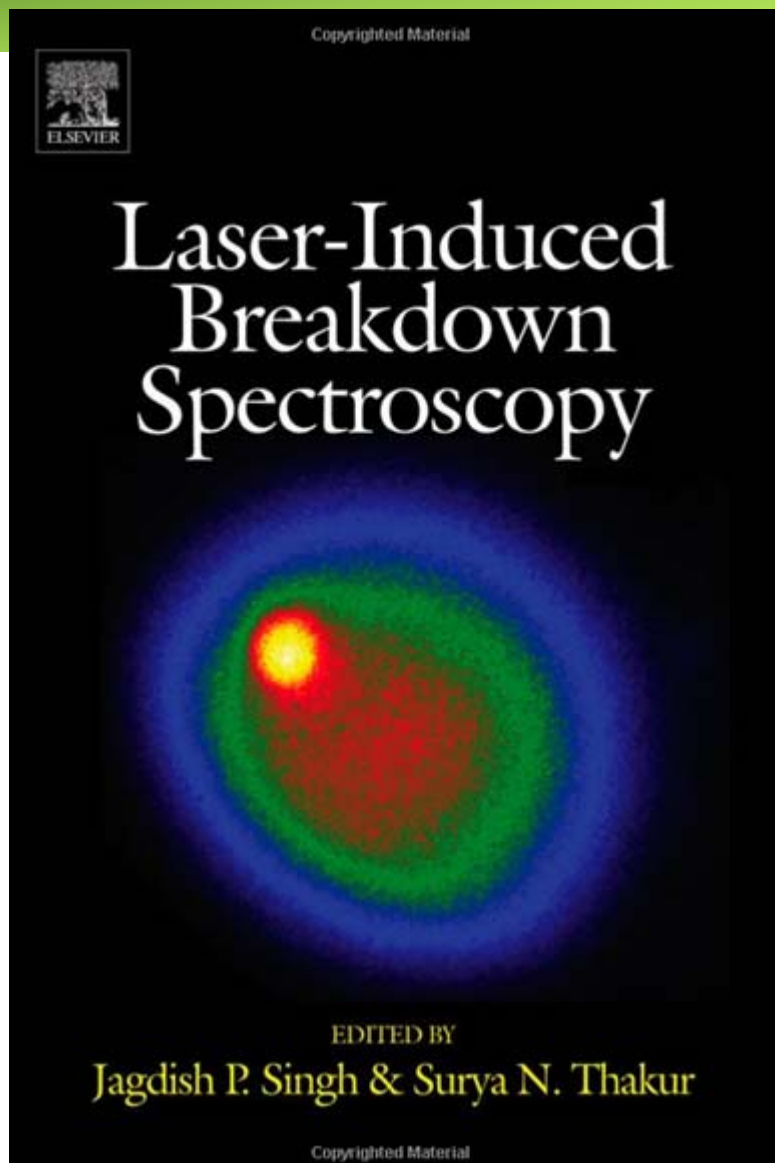
Leon J. Radziemski

Research Corporation, Tucson, AZ

*Hardback (ISBN-13: 9780470092996 | ISBN-10:
0470092998)*

Published July 2006 | 302 pages |

John Wiley & Sons, Ltd



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Edited by-

Jagdish P. Singh

Mississippi State University, Starkville, MS

Surya N. Thakur

Banaras Hindu University, Varanasi, India

Hardback (ISBN-13: 978-0444517340 | ISBN-10: 0444517340)

Published November 2007 | 454 pages |

Elsevier B.V.

Other Resources

- Progress in LIBS podcast (Dr. Andrzej Miziolek with “Spectroscopy”)
<http://spectroscopyonline.findanalytichem.com/spectroscopy/Podcast/PodcastProgress-in-LIBS/ArticleStandard/Article/detail/666685?contextCategoryId=49103>
- Video of real-time LIBS identification of explosive/biological threat (US Army Research Lab)
<http://www.arl.army.mil/www/default.cfm?Action=247&Page=462>
- LinkedIn group. Laser-Induced Breakdown Spectroscopy (LIBS)  This group is gathering all the peoples involved in the development and commercialization of the LIBS technique.
- <http://www.physics.wayne.edu/~srehse/> My website. I will post this talk there and have some other info about my projects.



Thank you for your attention and the invitation to join you today.

Questions?

